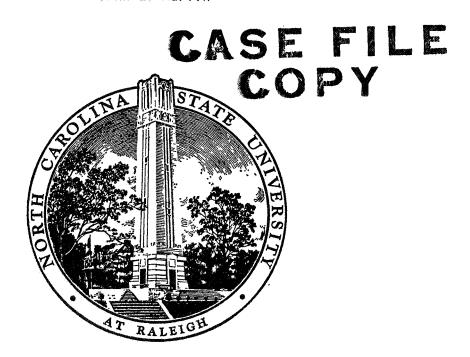
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INTRODUCTION

Increasingly, liquid hydrazine (N_2H_4) is coming to be considered as a convenient source of hydrogen rather than just as a rocket fuel. For such purposes, the hydrogen is usually obtained by passing the hydrazine through a heated catalytic bed. One convenient measure of the effectiveness of a catalytic decomposition device as a whole is to compare the quantity of hydrogen produced with the equilibrium concentration of the gaseous species N_2H_4 , NH_3 , N_2 , and H_2 which would exist at the temperature and pressure found in various parts of the device. Since such data do not seem to be readily available in the literature, calculations of the concentrations were carried out and are reported here.

The report begins with a review of the pertinant literature, citing experimental results or curve fits to the data where available. The basis for the computer program is then described. Following presentation of the results in both tabular and graphical forms is a comparison between the computed equilibrium concentrations and available experimental data.

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SURVEY OF THE LITERATURE

This section provides a brief survey of the available literature on the thermodynamic properties of hydrazine and its decomposition products as well as the decomposition reaction itself.

Several surveys of the thermodynamic properties of hydrogen, nitrogen, ammonia, and hydrazine are available in the literature. The first of these, by J. B. Evans and V. J. DeCarlo (1965) was intended as a bibliography for investigations into the chemistry of "cometary phenomena". A second bibliography of thermo-chemical data was published by M. K. Buresh, M. L. Reilly, G. T. Furukawa, and G. T. Armstrong (1965) of the National Bureau of Standards. This particular report listed the low-range temperature range data for the heat-capacities, enthalpy, and entropy of ammonia, carbon dioxide, and water. Both reports consisted of references to other investigations. Properties of gaseous nitrogen, hydrogen, ammonia, and hydrazine have been tabulated and are presented graphically in the N.A.S.A. Design Guide for Pressurized Gas Systems, Volume 11 (1966).

The data for hydrogen and nitrogen are well known and will therefore not be discussed further here.

J. L. Haws and D. G. Harden (1965) presented empirical equations for the thermodynamic properties of hydrazine, including vapor pressure, saturated liquid density, and the heat capacities. From the 'triple point' to the critical point, the vapor pressure is represented by

(1)
$$\ln P = 24.24 - (18,184.9/T) + 0.47629(\ln T) - 0.003863T + (1115.43(1190.08 - T)/1190.088) \ln(1190.8 - T)$$

for pressure in psia units and temperature in $^\circ R$ units. The error is relatively small. For the heat capacities, these authors give the following relations:

(2)
$$C_p = 0.137858 + 0.052715 \times 10^{-3} T - 0.119907 \times 10^{-6} T^2$$

(3)
$$C_v = 0.075903 + 0.52715 \times 10^{-3} \text{T} - 0.119907 \times 10^{-3} \text{T}^3$$

An equation of state for gaseous hydrazine was presented in the form

(4)
$$P = RT/(V - b) + \frac{A_2 + B_2T + C_2e^{-xT/T_c}}{(V - b)^2} + \frac{A_3 + B_3T + C_3e^{-xT/T_c}}{(V - b)^3}$$

$$+ \frac{A_4}{(V - b)^4} + \frac{B_5T}{(V - b)^5}$$

where R = 10.7315/mole wt

 $b = 0.0299451 \text{ ft}^3/1\text{bm}$

x = 8.0000

 $T_{c} = 1175.69^{\circ}R$

 $A_2 = -32.21898$

 $B_2 = 0.00294549$

 $C_2 = -1171.2876$

 $A_3 = 0.869588$

 $B_3 = 0.16838487 \times 10^{-3}$

 $C_3 = 46.250721$

 $A_4 = -0.202100 \times 10^{-1}$

 $B_5 = 0.12899897 \times 10^{-6}$

The equations were used in computing property values for a temperature range from 32 °F to 1400°F, and for pressures from 1.0 psia to 8000 psia. It should be noted that the equations were curve-fitted to match experimental data provided by other investigators.

M.S. Jhon, J. Grosh, and H. Eyring (1967), using the significant structure theory of liquids, computed the thermodynamic properties of liquid hydrazine (and other chemicals). The computational results were claimed to be in good agreement with experimentally-determined property data. These properties were I) the vapor pressure, 2) the molar volume, 3) the entropy, 4) the critical temperature and pressure, 5) the constant volume and constant pressure heat capacities, 6) the coefficient of thermal expansion, and 7) the compressibility. The surface tension and the dielectric constant were computed by other methods which are related to the significant structure theory of liquids.

The "significant structure" theory for liquids is actually a model approach for describing the mechanical and thermodynamic properties of liquids.

Let V be the mole volume of the liquid state, V_S be the mole volume of the solid state at the melting point, and V_O be the mole volume of the "solidlike structure" in its molten state. The "excess volume" of the liquid is defined as $V-V_S$. This volume is due to expansion of the chemical substance upon melting and entering the liquid phase. The liquid is considered to have molecular "vacancies" in its interior, and the total volume of these vacancies is the "excess volume". This volume permits molecules in the liquid state to have translational degrees of freedom.

If a molecule is completely surrounded by other molecules, then it can only oscillate. The oscillation frequency will be nearly the same as if the molecule were in a solid state. If the molecule is adjacent to a "vacancy", then the molecule will have translational degrees of freedom, and will behave as if it were in a gaseous phase.

The partition function used in the theory is formed from the multiplication of the partition functions for each type of structure. A factorial term is included to account for the lack of distinction between the 'gas' molecules.

The liquid state partition function was written in the form

(5)
$$f = \frac{\begin{bmatrix} f_{\text{solid}} & f_{\text{deg}} \end{bmatrix} & \begin{bmatrix} f_{\text{gas}} \end{bmatrix}}{\begin{bmatrix} N(V - V_0) \end{bmatrix}}$$
where $f_{()}$ are the partition functions for solid and the solid solid and the solid soli

where $f_{\mbox{\scriptsize (}}$ are the partition functions for solid and gas states, and the molecular degeneracy of position.

The Helmholtz free energy function was expressed as

(6)
$$A = -kT \ln (f).$$

Using the above equation, the vapor pressure was computed in the following manner. For a constant temperature, the Helmholtz free energy was plotted as a function of volume. Through points common to the liquid and vapor states, a tangent was drawn, whose slope gives the vapor pressure. The values of volume for the two tangency points give the values for the liquid and vapor molar volumes.

The entropy of vaporization was computed from the relation

(7)
$$S = -\left[\frac{\partial A}{\partial T}\right]V = \left[\frac{\partial}{\partial T}(kT \ln f)\right]V$$

and the vaporization entropy was given by S_g – S_l , where lg' and ll' indicate vapor and liquid states respectively. The constant-volume heat capacity is given by

(8)
$$C_{V} = \frac{\partial}{\partial T} \left[kT^{2} \frac{\partial \ln f}{\partial T} \right]_{V}$$

and the constant-pressure heat-capacity is given by

$$C_{p} = C_{v} + TV \frac{\alpha^{2}}{\beta}$$

where

(10)
$$\alpha = \frac{1}{v} \left[\frac{\partial V}{\partial T} \right]_{P} = -\frac{1}{v} \left[\frac{\partial^{2} A}{\partial V \partial T} \right]_{TV} / \left[\frac{\partial^{2} A}{\partial V^{2}} \right]_{T}$$

which is the thermal expansion coefficient, and

(11)
$$\beta = -\frac{1}{V} \left[\frac{\partial V}{\partial P} \right]_{T} = \frac{1}{V} V \left[\frac{\partial^{2} A}{\partial V^{2}} \right]_{T}$$

which is the compressibility. The critical temperature, pressure, and volume were computed by an iteration method using the conditions

$$\frac{\partial P}{\partial V} = 0$$

$$\frac{\partial^2 P}{\partial V^2} \bigg|_{T} = 0$$

(14)
$$P = -\left[\frac{\partial A}{\partial V}\right] T$$

The surface tension was computed by an iteration technique which accounted for polar liquid orientation effects. The dielectric constant was expressed as a function of the index of refraction and the dipole moment.

A comparison of the results of Haws and Harden, with those of Jhon, Grosh, and Eyring, for vapor pressure and constant pressure heat capacity show discrepancies. Only for the saturated liquid density is there reasonable agreement. The following table compares the results of the two computations for the vapor pressure of N_2H_4 .

Table I. Comparison of computed vapor pressures.

T, °K	P, atm(Haws and Harden)	P, atm(Jhon <u>et al.</u>)
273	0.003539	0
293	0.0135658	0.01
313	0.04322	0.04
333	0.11850	0.20
343	0.18696	0.27
1 .		

No comparison is possible for temperatures above 400 °K.

A comparison of the constant-pressure heat capacities gives the following:

Table 2. Comparison of computed constant-pressure heat capacities.

T, °K	c _p , cal/g-mole-°K (Haws and Harden)	c _p , cal/g-mole-°K (Jhon <u>et al</u> .)
273	11.7792	20.84
298	12.3854	21.27
300	12.4064	21.35
400	14.5728	25.575

The comparison of results for both vapor pressure and constant-pressure heat-capacities show that Jhon, Grosh, and Eyring consistently obtained higher values at elevated temperatures than did Haws and Harden.

By contrast, the results obtained in each study for the saturated liquid density were in good agreement.

Table 3. Comparison of computed specific volumes.

т , ° К	V, cc/mole (H&H)	V, cc/mole (J,G,&E)		
273	31.261	31.327		
293	31.7269	31.67		
308	32.1294	31.984		
400	35.86127	34.016 @ 387.3 (bp)		
1				

If the results of Haws and Harden, and those of Jhon, Grosh, and Eyring are compared with the properties tabulated in the N.A.S.A. Design Guide, the results of Haws and Harden are found to compare favorably with those listed in the Design Guide. For the purposes of the present study, then, the equations of Haws and Harden will be used to provide standards against which to compare the results of numerical computations made using the N.A.S.A. equilibrium composition analysis computer program. These computations are described in detail in subsequent sections of the present work. For the moment it is sufficient to note that in these computations each specie is assumed to be a perfect gas.

Of the major products of hydrazine decomposition, the only one which can liquify in the range of temperatures and pressures of present interest is ammonia. Since a gas is usually quite imperfect thermally near saturation its exact thermodynamic properties under these conditions are of considerable interest to one wishing to refine equilibrium composition calculations.

The thermodynamic properties of ammonia have been tabulated and graphically presented in the N.A.S.A. Design Guide (1966). Specific properties of interest, which are listed, are I) the vapor pressure of liquid ammonia, and 2) the density of saturated ammonia vapor, 3) the density of liquid ammonia as a function of both pressure and temperature, 4) the heat capacities, and 5) the heat of vaporization.

The heat of formation of gaseous ammonia is -11.02 kcal/mole (Penner, 1968). Equations of state which may be used for ammonia are 1) the Berthelot equation, 2) van der Waals, 3) Beattie-Bridgeman, and 4) the Dieterici equation. The equations given below are taken from Penner (1968). The Berthelot equation is

(15)
$$V = \frac{RT}{P_{c}} + \frac{9RT_{c}}{128p_{c}} \left[1 - \frac{6T_{c}^{2}}{T^{2}} \right]$$

where p_{C} and T_{C} are respectively the critical pressure and temperature. The gas constant R for ammonia is obtainable, for any set of units, from the N.A.S.A. Design Guide.

The van der Waals equation would give

(16)
$$\left[p + \frac{a}{\sqrt{2}}\right] \left[V - b\right] = RT$$

where a = $3p_cV_c^2$ and b = $V_c/3$.

The Beattie=Bridgeman equation for gaseous ammonia is

(17)
$$p = RT \left[I - \frac{c}{VT^{3}} \frac{I}{V^{2}} \left[V + B_{o} - \frac{bB_{o}}{V} \right] - \frac{A_{o}}{V^{2}} - \frac{aA_{o}}{V^{3}} \right]$$

where $A_0 = 2.3930$

 $B_0 = 0.03415$

a = 0.17031

b = 0.019112

 $c = 476.87 \times 10^{-4}$

for units in atmospheres, liters per mole, and °K.

The Dieterici equation, which is very accurate near critical points, is

(18)
$$\left[\frac{a}{Pe^{VRT}} \right] (V - b) = RT$$
 where $a = 7.39p_cV_c^2$ and $b = V_c/2$.

The values for the constants in these equations are from Hirschfelder, Curtiss, and Bird (1954).

The remainder of this review is concerned with papers treating the decomposition of hydrazine.

K. W. Michel and H. Gg. Wagner (1965) studied the thermal decomposition of $\rm N_2H_4$ diluted with argon and/or helium, behind shock waves, in the temperature range II00 °K to I600 °K, for pressures from 26.65 atm to 80.0 atm. For temperatures between II00 °K and I200 °K, decomposition did not start until an 'induction period' had passed. For temperatures greater than I300 °K, the decomposition reaction begins by following a first-order reaction-rate law, and then changes to successive reactions of varying order.

A reaction of order one follows a rate law such as

$$-\frac{dc}{dt} = k_1 c$$

where c = specie concentration, k_1 = the rate constant, and t = time. Reactions may have an order greater than one. An example of such a reaction is

$$(20) \qquad A + B \longrightarrow X + Y.$$

Let c_A be the concentration of specie A and c_B be the concentration of specie B. Let the superscript o indicate the value of a quantity at time t=0. The reaction rate equations are

(21)
$$-\frac{dc_A}{dt} = -\frac{dc_B}{dt} = k_2 c_A c_B.$$

Define
$$x = c_A^O - c_A = c_B^O - c_B$$

(22)
$$\frac{dx}{dt} = k_2(c_A^0 - x)(c_B^0 - x).$$

If, as a special case, $c_A^o = c_B^o$, then

$$\frac{dx}{dt} = k_2(c_A^o - x)^2$$

$$(24) \qquad -\frac{dc_A}{dt} = k_2 c_A^2.$$

Many reactions are not as simple as those given above, and the reaction order may be non-integer in value.

The results of Michel and Wagner's investigation confirmed the "straight-chain" model for the decomposition reaction which had been proposed by G. K. Adams and G. W. Stock (1953). Reaction models for the hydrazine break-down are discussed later. At higher temperatures, the reaction rates were found to be partially dependent on the total gas density. For a density of 7.5 \times 10⁻⁵ moles/cm³, (26.65 atm),

(25)
$$\log k = 12.8 - ((52,000 \text{ cal/mole})/2.3RT)$$

and for a density of 2.5×10^{-5} moles/cm³, (80.0 atm)

(26)
$$\log k = 12.0 - ((48,000 \text{ cal/mole})/2.3RT)$$
.

The main objective of the investigation was to determine the effects of small amounts of oxygen on the decomposition rate. It was found that small amounts (i.e. 0.2% of the mixture) do not affect the reaction. For temperatures above 1400 °K, heavy concentrations of oxygen, (e.g. $[0_2]/[N_2H_4] = 10$) still do not affect the decomposition reaction.

The amount of ammonia formed during the decomposition was also determined. It was found that at 1100 °K the ratio of NH $_3$ formed per amount of N $_2$ H $_4$ decomposed is equal to one, at 1600 °K equal to 0.5, and at 2000 °K equal to zero. The radical NH appeared for temperatures above 1400 °K.

R. Roback (1965) made an investigation into the thermodynamic properties and equilibrium compositions of the products of dissociation of hydrogen, ammonia, and other materials, such as graphite, molybdenum, tungsten, and the carbides, nitrides, and oxides of titanium and zirconium. The temperature range was from 1800 °R (1000 °K) to 10,000 °R (5550 °K) for a pressure of 1000 atm. The materials considered were for use in controlling the radiation heat transfer to the moderator walls of gaseous nuclear-engined missles. The computations were carried out by assuming that all species obeyed the ideal gas law. Viscosities and thermal conductivities were obtained using data valid at one atm pressure, and a pressure correction determined from the Enskog theory for dense gases. Atomic and molecular species were considered

in the computation, but ionic species were considered to be of negligible importance below a temperature of 10,000 °R (5550 °K). The calculations indicated that hydrogen dissociation is important for temperatures above 7000 °R (3880 °K). By contrast, concentrations of ammonia are almost completely decomposed into $\rm H_2$ and $\rm N_2$ at 2000 °R (III0 °K). Densities and enthalpies were computed for hydrogen, ammonia, and other gases, and presented as functions of temperature.

I. J. Eberstein and I. Glassman (1965) experimentally studied the decomposition of low concentrations of hydrazine for a temperature range of 750 °K to 1000 °K. The apparatus used was an adiabatic flow reactor with a driving pressure of 136 atm. Both the rate constants and reaction stoichiometry were determined. For hydrazine break-down inside a 3-inch duct, the reaction rate constant was determined to be

(27)
$$k = 10^{10.33} \exp(-36.170/RT) \sec^{-1}$$

The over-all reaction, of an order close to one, was given as

(28)
$$N_2H_4 \rightarrow 0.9NH_3 + 0.5N_2 + 0.6H_2$$
.

The stoichiometry of the reaction was determined by experiment.

Hydrazine-water mixture decomposition was also studied. For mixtures where the water was more then one percent of the total volume, the rate of decomposition was decreased by a factor of ten. The reaction rate for these mixtures was found to be independent of the amount of water present in the mixture, provided that the amount of water was more than one percent of the total volume. It was concluded that water must suppress one of the decomposition reaction steps.

Along with hydrazine, both unsymmetrical dimethyl hydrazine and monomethyl hydrazine were considered. Of the three, anhydrous hydrazine, N_2H_4 , had the slowest decomposition rate.

E. J. Bair et al. (1966) used the explosive decomposition of hydrazine as part of their experimental program in the development of a high-resolution spectrometer and the spectroscopic procedures in examining explosive reactions. The reaction was described as a series of "free radical chain processes". One of the first products to be formed was NH2, but this product dissappeared as the reaction continued, forming NH. The NH radical in turn also disappeared. Both radicals were observed spectrometrically. When the NH radical vanished, a spectrum of lines in the 0.600 μ region was found which was thought to arise from the formation of N2H3, termed the "hydrazil radical".

- A. S. Sokolik, V. P. Korpov, and E. S. Semenov (1967) used hydrazine decomposition flames in their study of turbulent burning rates and flame velocities. The decomposition process was considered to be a "straight-chain mechanism." The breaking of the reaction chain, however, is dependent on pressure. The reaction order is also dependent on reaction pressure. At low pressures, the order is approximatedly 1.0, and at high pressures the order is 1.5. It was noted that hydrazine "decomposes extremely readily," requiring special care in both the handling and experimental procedures.
- R. F. Sawyer and I. Glassman (1967) have examined the reactions between hydrazine and nitrogen dioxide, nitric oxide, oxygen, and oxygen-nitric oxide, by using an adiabatic flow reactor. For each reaction, the investigators determined reaction rates, activation energies for the over-all reaction, and the reaction orders. Reacting with oxygen, hydrazine forms water and nitrogen according to the reaction

(29)
$$N_2H_4 + O_2 \longrightarrow 2H_2O + N_2$$

For a total pressure of 1.0 atm and temperatures between 950 °K and 1010 °K, the hydrazine/oxygen reaction rate law is expressed as

$$\frac{d[N_2H_4]}{dt} = -k[N_2H_4]$$
with

(31)
$$k = 10^{9.91} \exp(-37,200/RT) \sec^{-1}$$

and a reaction order of one. The activation energy was 37.2 kcal/mole.

The reaction rate is dependent on the concentration (to power one) and independent of the concentration of oxygen. Oxidation was shown to be possible whether or not the hydrazine was decomposed prior to the reaction. The production of water by the reaction was believed to be responsible for slowing the rate of decomposition.

D. I. Maclean and H. Gg. Wagner (1967) studied the decomposition-flame structure for both ammonia-oxygen and hydrazine vapor, determining experimentally temperature and specie-concentration profiles. The initial temperatures of the hydrazine vapor were between 70 °C and 90 °C, for pressures of 0.01842 atm and 0.0342 atm. It should be noted that the only initial substance in combustion was the pure hydrazine, with no other chemical species present. The reaction, after ignition, proceeds without aid of any oxidizers, and is exothermic. For the hydrazine-vapor flame, a maximum temperature of 1350 °K was found by means of a thermocouple. The investigators believed the temperature reading to be too low, as the adiabatic flame temperature should be near 1700 °K.

It was found by experiment that increasing the flame speed decreased the rate of decomposition of the first few millimeters above the burner. A major factor in the rate of decomposition is pressure. By doubling the pressure, nearly twenty mole percent of the hydrazine was broken down within the first millimeter above the burner.

The concentrations of $\rm H_2$, $\rm N_2$, and $\rm NH_3$, which are the products of the reaction, and the rate of the reaction, are not affected by changes in the initial temperature of the hydrazine vapor. The maximum temperature is not affected by changes in pressure.

E. F. Logan and J. M. Marchello (1968) decomposed hydrazine by means of an alternating-current glow-discharge flow reactor. Using several flow rates and electric currents, light emmission spectra were measured from 0.300 μ to 0.600 μ . The most intense spectrum was produced by molecular nitrogen. However, atomic nitrogen spectra were not observed. Both atomic and molecular hydrogen were observed. The emmission bands for ammonia were not found. The radicals NH, NH2, were found but OH was not. The gaseous hydrazine vapor pressure was 1.91 \times 10 $^{-2}$ atm.

It was found that 65% of the hydrazine is decomposed very rapidly, with the rest of the reactant decomposing at a slower rate.

The decomposition of hydrazine has also been used in the study of the dissociation energy of the N-H and N-N bonds.

I. P. Fisher (1965a, b), with G. A. Heath, experimentally determined the value for the dissociation energy of the N-H bond. The value was found by electron-impact measurements, and by using the relation

(32)
$$D(H - N_2H_3) = A(N_2H_3^+) - I(N_2H_3)$$

where A is the appearance potential of $N_2H_3^+$, I is the ionization potential of N_2H_3 , and D is the dissociation energy of the N-H bond. The values which were found were

(33)
$$A(N_2H_3^+) = 11.3 \pm 0.1 \text{ ev}$$

(34)
$$I(N_2H_3) = 7.88 \text{ eV}$$

(35)
$$D(H-N_2H_3) = 78.0 \text{ kcal/mole}.$$

Values for bond strengths, and heats of formation and ionization potentials of several chemical species resulting from hydrazine decomposition were also determined (Fisher, 1965b).

Another study of bond dissociation energies, for the N-N bond, was made by H. A. Olschewski, J. Troe, and H. Gg. Wagner (1966). It was shown that the decomposition activation energy must be measured for decomposition at high pressures. A shock tube of 70 mm diameter was constructed and used in their experimental program, with driving pressures up to 200 atm being applied.

For decomposition at high pressures, and densities greater than 5×10^{-4} moles/cm³, the experimental results provided the temperature dependence of the reaction rates. The reaction rates were then used in determining the unimolecular decomposition activation energy, which was found to be $55^{\pm}2$ kcal/mole. The reaction rate for decomposition at high pressures was determined to be

(36)
$$k = 10^{13.6} \exp(-55,000/RT) \sec^{-1}$$
.

This rate constant is for the over-all reaction, and not to be confused with the dissociation energy of the N-H bond.

Two papers were found which discussed the decomposition of hydrazine by catalytic means.

A. S. Keston (1967) carried out a theoretical study of catalytic reactors designed for hydrazine decomposition.* Transient and steady-state models were developed for temperature and specie-concentration profiles within the reactor, as functions of time and position. Both thermal and catalytic-induced decomposition were considered, with the inclusion of the effects of heat and mass transfer between the free-flowing gas, and gas contained within the catalyst material. The results of the computations were compared with data from hydrazine-fueled engine experiments where the Shell 405 catalyst was used, with very good agreement between the theoretical and experimental results. These models provide an indirect means by which the decomposition of hydrazine and ammonia may be computed.

A summary of the results of the steady-state reaction calculations of Keston is given on the following page. The values are for an axial distance of 0.25 feet from the inlet end of the catalytic reactor.

^{*} Temperatures varied from 1000 °K to 1330 °K and pressures varied from 7.6 atm to 71 atm.

Table 4. Summary of steady-state reaction computations of Keston

Р	T °K	Mole Flow Rate	Mole Fractions			Fractional
atm.	°K	lbm/ft ² -sec	H ₂	N ₂	NH ₃	NH ₃ dissoc.
7.59	1000	1.51	0.55 (0.59)	0.31	0.14 (0.10)	0.62
14.8	1020 (1000)	1.52	0.57 (0.59)	0.32 (0.33)	0.10 (0.07)	0.71
32.6	1054 (1076)	3.12	0.55 (0.56)	0.31	0.14	0.64
66.2	1120	6.29	0.53 (0.51)	0.31 (0.32)	0.16 (0.17)	0.57
71.0	1054 (1076)	2.43	0.57 (0.54)	0.32	0.12 (0.15)	0.69

The numbers in (-) represent experimental data differing from the theoretical computations.

The mole fractions of $\rm H_2$ and $\rm N_2$ produced by the reaction are nearly independent of variations in pressure and temperature. The mole fraction of hydrogen produced varied between 0.53 and 0.57, and the mole fraction of nitrogen produced varied from 0.31 to 0.32. Ammonia mole-fractions varied between 0.10 and 0.16. It was found that in the reaction, mole-fractions of ammonia reached a maximum with respect to axial distance in the reactor (or equivalently, time), and then declined in value. Between 57% and 71% of the NH $_3$ produced in the reaction decomposes to nitrogen and hydrogen.

It may be noted that for the maximum mass flow rate, hydrogen production was at a maximum, as was the fractional NH₃ dissociation. However, the temperature of the reaction reached a maximum for lower mass-flow rates.

At either 14.8 atm or 71.0 atm, there was maximum output of hydrogen, chiefly due to the maximum dissociation of ammonia.

A second study of catalyst-induced hydrazine break-down was made by C. F. Sayer (1970) who also used the Shell 405 catalyst. The purpose of the study was to examine the chemical kinetics and mechanisms involved in the break-down of liquid hydrazine in the presence of the catalyst.

Using the Shell 405 catalyst, the break-down resulted in a gas mixture which was reported to be 99.4% $\rm N_2$ and 0.6% $\rm N_2H_4$. The vanishing of the hydrogen was attributed to "the conversion of the hydrazine molecule into the nydrazinium ion." If a rhodium catalyst was used instead of the Shell 405 catalyst, then the resulting gas mixture was nearly 50% $\rm N_2$ and 50% $\rm H_2$. It may be noted that a main component of the Shell 405 catalyst is iridium. The reaction with the Shell 405 catalyst was expressed as

$$(37) 3N_2H_4 \longrightarrow N_2 + 4NH_3.$$

For the rhodium catalyst, the reaction was written as

$$(38) 2N_2H_4 \longrightarrow N_2 + H_2 + 2NH_3.$$

For the production of hydrogen from hydrazine, the rhodium catalyst is obviously to be preferred.

Two references attest to the volatility of hydrazine. G. P. Sutton (1963) noted that hydrazine may spontaneously ignite with either nitric acid or hydrogen peroxide. It was also noted that hydrazine vapor may form explosive mixtures with air.

Sutton very briefly discussed the reaction of hydrazine with catalysts. The products of the decomposition reaction, the exact composition of the product mixture, and the pressure will depend on the type of catalyst used.

If impurities are present in liquid hydrazine, or in contact with hydrazine vapor, then the hydrazine will readily decompose exothermically.

The reaction was stopped when the total pressure in the experimental apparatus reached 1.025 atm.

E. J. Bowen and A. W. Birley (1951) found by experiment that for very low pressures, and temperatures of the order of 100 °C, that mixtures of hydrogen peroxide and hydrazine vapors may co-exist for a brief period of time without reaction. However, decomposition will then take place.

A comparison among the results for hydrazine decomposition shows fair agreement between the investigations of Michel and Wagner (1965), Eberstein and Glassman (1965), Sawyer and Glassman (1966), and Maclean and Wagner (1966). Without the presence of a catalyst, the stoichiometry of the decomposition reaction is generally

(39)
$$N_2H_4 \longrightarrow NH_3 + \frac{1}{2}H_2 + \frac{1}{2}N_2$$
.

By inducing decomposition with a glow-discharge at a pressure of 1.91 x 10^{-2} atm, Logan and Marchello (1968) obtained a gas mixture consisting of 46% $\rm H_2$, 28% $\rm N_2$, and 19% NH $_3$ for 85% hydrazine break-down.

C. F. Sayer (1970) found that catalyst-induced reactions give results different from the previous investigations mentioned above. Using the Shell 405 catalyst, one mole of hydrazine results in the formation of 1.33 moles of ammonia. The final gas composition is however 99.4% N2, and 0.6% N2H4. If a rhodium catalyst is used, one mole of hydrazine decomposes into one mole of ammonia, with a final gas composition of 50% N2, and 50% H2. The chemical reactions were stopped when total system pressures exceeded 1.0 atm.

Michel and Wagner 26.26 atm to 80.0 atm

Eberstein and Glassman 136 atm

Sawyer and Glassman I atm

Maclean and Wagner 0.0184 atm to 0.0342 atm

Reaction pressures varied considerably. The pressures in the experiments were

A. S. Keston (1967), by theoretical means, found that for elevated pressures, the results of decomposition of hydrazine are mostly $\rm H_2$ and $\rm N_2$. As an example, for a pressure of 7.6 atm, and a temperature of 1000 °K the gas mixture consisted of approximately 54% $\rm H_2$. 31% $\rm N_2$, and 15% NH₃. The maximum temperature was 1120 °K, and the maximum pressure was 71.0 atm.

For relatively low pressures, and temperatures of the order of 1000 °K, the present numerical investigation shows that the products of the decomposition reaction are almost entirely hydrogen and nitrogen, with twice as many moles of hydrogen being produced as nitrogen.

In comparing these results with those of previous investigations, the agreement is not unreasonable. In previous studies, for very low pressures, one mole of hydrazine decomposed to approximately one mole of ammonia, and a half-mole each of hydrogen and nitrogen. If a rhodium catalyst is used, the products of the reaction are equal mole numbers of hydrogen and nitrogen, with no ammonia left at the end of the reaction.

For elevated pressures (Keston, 1967), hydrogen is the major product (54%) with lesser amounts of nitrogen (31%) and ammonia (15%). It may be seen that as pressure increases, for temperatures in the neighborhood of 1000 °K, the ammonia production declines, and hydrogen and nitrogen are produced in amounts which appear to approach a ratio of two to one. The present computational prediction is that for pressures in the neighborhood of 7.6 atmospheres and temperatures of about 1000 °K, the mole-fractions of the decomposition products are 34% nitrogen and 66% hydrogen.

Several reaction mechanisms have been postulated for the break-down of hydrazine.

Michel and Wagner (1964) have noted that the first two reaction steps have been theoretically determined and experimentally verified. These steps are

(41)
$$NH_2 + N_2H_4 \rightarrow N_2H_3 + NH_3$$

For low temperatures, the next steps would then be

followed by

$$(43) \qquad \qquad H + N_2H_4 \longrightarrow NH_3 + NH_2$$

or

(44)
$$H + N_2H_4 \longrightarrow H_2 + N_2H_3$$

$$(45) \qquad NH_2 + HNNH \longrightarrow NH_3 + N_2 + H$$

$$(46) 2NH2 \rightarrow NH3 + NH$$

$$(47) \qquad M + NH_2 + H \longrightarrow NH_3 + M .$$

At temperatures of 1600 °K, the step

$$2NH_2 \longrightarrow HNNH + H_2$$

preceeds reaction (6). For temperatures in the neighborhood of 2000 $^{\rm o}$ K, step (6) could be replaced by the reaction

$$(49) \qquad \qquad +NNH \longrightarrow NNH + H.$$

Eberstein and Glassman (1965) considered two possibile reaction mechanisms. The simple mechanism was similar to the postulated mechanism of Michel and Wagner, except for the last step. Michel and Wagner assumed the reaction

$$(50) \qquad \qquad NH_2 + HNNH \longrightarrow NH_3 + N_2 + H$$

whereas Eberstein and Glassman consider the reaction to be

(51)
$$N_2H_3 + NH_2 \longrightarrow NH_3 + N_2 + H_2$$
.

The assumed mechanism includes branching reactions. By numerical techniques, the investigators found that the reaction has four parts. In the first part, the formation of free radicals takes place. The reaction steps which follow can be characterized by "steady-state" rate constants. The third part of the reaction was characterized by decreasing first-order rate constants. In the last part, the concentrations of free-radicals and the over-all reaction rate rapidly declined. When compared with experimentally-measured stoichiometry and reaction-rates, the reaction mechanism was found to be "in good quantitive agreement."

The reaction mechanism which was postulated is tabulated as follows.

Initiation:

$$(52) N2H4 + X \longrightarrow 2NH2 + X.$$

Propagation:

$$(53)$$
 $N_2H_4 + NH_2 \longrightarrow NH_3 + N_2H_3$

(54)
$$N_2H_3 + X \longrightarrow N_2 + H_2 + H + X$$

$$(55) N2H4 + H \longrightarrow NH3 + NH2.$$

Branching:

(56)
$$N_2H_3 + X \longrightarrow NH + NH_2 + X$$

$$(57) N2H4 + NH \longrightarrow NH2 + N2H3,$$

Termination:

(58)
$$NH_2 + N_2H_3 \longrightarrow NH_3 + N_2 + H_2$$

(59)
$$N_2H_3 + N_2H_3 \longrightarrow 2NH_3 + N_2$$

$$(60) N2H3 + H \longrightarrow N2 + 2H2$$

$$(61) \qquad \qquad NH_2 + NH_2 \longrightarrow N_2H_4 \bullet$$

Sawyer and Glassman (1966) have tabulated the assumed mechanisms for reactions of hydrazine with nitrogen dioxide, nitric oxide, and oxygen. In the reactions given, the authors assumed that branching occurs before propagation.

Sokolik, Korpov, and Semenov (1967) have noted that the "chain" mechanism of Adams and Stocks is usually accepted. The steps for this mechanism are

$$(62) N2H4 \longrightarrow 2NH2.$$

Continuation: (propagation)

(63)
$$N_2H_4 + NH_2 \longrightarrow NH_3 + N_2H_3$$

$$(64) \qquad \qquad N_2H_3 \longrightarrow N_2 + H_2 + H_3$$

$$(65) \qquad \qquad H + N_2H_4 \longrightarrow NH_3 + NH_2.$$

The chain breaking is dependent on the reaction pressure. For low pressures, the branching and termination steps are

(66)
$$H + N_2H_3 \longrightarrow N_2 + 2H_2$$

(67)
$$NH_2 + N_2H_3 \longrightarrow NH_3 + N_2 + H_2$$

For high pressure reactions, the branching and terminal steps are

$$(68) \qquad \qquad H + NH_2 + X \longrightarrow NH_3 + X$$

$$(69) \qquad \qquad H + H + X \longrightarrow H_2 + X.$$

Logan and Marchello (1968), discussing the glow-discharge break-down of hydrazine, gave the following reaction steps as the mechanism of electron-induced decomposition. The first step is caused by the electron interaction

(70)
$$N_2H_4 + e^* \longrightarrow N_2H_4^* + e$$

The (*) indicates an activation state. The authors' spectrometer measurements indicated that $N_2H_A^*$ breaks down in two simultaneous reactions.

(71)
$$N_2H_4^* \longrightarrow N_2H_3 + H$$

$$(72) N2H4* \longrightarrow NH3 + NH$$

The reactions which then follow are

(73)
$$H + N_2H_3 \longrightarrow H_2 + N_2H_2$$

(74)
$$NH + N_2H_4 \longrightarrow NH_3 + N_2H_2 \bullet$$

The diazene (or diimide) decomposes as

$$(75) \qquad \qquad N_2 H_2 \longrightarrow N_2 + H_2.$$

In all of the assumed reaction mechanisms, N_1H_4 always participates in a later reaction step, as well as in the initial step. There is disagreement as to the exact steps of the reaction after the initial and propagation reactions.

Sayer (1970) formulated a hydrazine break-down mechanism involving the use of catalysts, to explain the results of his experiments. It was postulated that hydrazine molecules adsorb on the catalyst surface, and react with other hydrazine molecules which are "in solution" to form two "cis-hydrogen" atoms. Additional products of this reaction are two molecules of ammonia, with cis-diimide adsorbed on the catalyst surface. Two adjacent molecules of diimide then react to form nitrogen and hydrazine. The nitrogen desorbs from the catalyst surface, and is replaced by another hydrazine molecule.

H. S. Maddix (1965) discussed the break-down of ammonia in his study of gaseous discharge phenomena. The ammonia decomposition was induced by an electrical discharge in a tube. The observed reaction was

$$(76) 2NH_3 \longrightarrow N_2 + 3H_2.$$

As the amount of ammonia decreased, the nitrogen concentrations steadily increased. The hydrogen concentrations first rapidly increased, the decreased, and then increased again. The changes in the amount of hydrogen were due to chemical adsorption of the hydrogen with the tube walls.

A second discussion of the decomposition of ammonia was provided by Roback (1965). It was noted that above 2000 °R, ammonia is almost entirely dissociated, the products being $\rm H_2$ and $\rm N_2$. Above 4000 °R, H, NH, and NH₂ concentrations become appreciable. For femperatures above 7000 °R, $\rm N_2$ begins to decompose to single atoms of nitrogen.

Michel and Wagner (1965) included a study of the pyrolysis of NH_2 , which showed that ammonia decomposition is a first-order reaction. For temperatures between 2100 $^{\circ}$ K and 2900 $^{\circ}$ K, and with a 'carrier' gas of argon, the reaction rate was

(77)
$$k = 10^{15.64} \exp[-(79.5 \pm 2.5 \text{ kcal/mole})/\text{RT}] \text{ cm}^3/\text{mole-sec}$$

For temperatures greater than 1400 $^{\circ}$ K, NH radicals will be present in the decomposition-products mixture.

As has been shown, values of thermodynamic properties of hydrogen, nitrogen, ammonia, and hydrazine are readily available. The properties of hydrazine have been determined from both experimental studies and theoretical computations. Empirical equations for the properties of hydrazine have been presented, and equations of state for both gaseous hydrazine and ammonia have been listed.

In the foregoing, previous studies of hydrazine decomposition have been summarized and discussed. The studies have shown that in general one mole of hydrazine decomposes into one mole of ammonia, one-half mole of hydrogen, and one-half mole of nitrogen. The exact stoichiometry of the reactions depends on both the temperatures and pressure of the reaction. The use of a catalyst can lead to increased production of both hydrogen and nitrogen. However, if the incorrect catalyst is used, very little hydrogen will result from the hydrazine decomposition.

The results of the previous investigations and present numerical study have, where possible, been compared.

As has been shown, the decomposition mechanism is quite complex, and there is still disagreement among researchers studying the hydrazine breakdown as to the exact steps in the reaction. Several of the investigations have included studies of ammonia decomposition, and these have been summarized.

It is clear that the properties of hydrazine are not as well determined as one would wish. Although previous studies are not in exact agreement, there is qualitative agreement among them. Much more experimental work, however, will have to be performed to determine the exact kinetics of the decomposition reaction mechanism.

THE COMPUTATIONAL METHOD

The computer program used in the present calculation of hydrazine decomposition was written by S. Gordon and B. J. McBride (1971), and listed in N.A.S.A. SP-273. The program was designed to compute the equilibrium composition of reacting chemical mixtures for given thermodynamic states, i.e. specified (T,P), (H,P), (S,P), (T,V), (U,V), or (S,V), where:

T = temperature, °K

 $P = pressure, newton/meter^2$

 $V = \text{specific volume, meter}^3/\text{kg}$

H = enthalpy, joules

S = entropy, joules/°K

U = internal energy, joules.

The chemical species may be either in a gaseous phase and/or in a condensed phase.

The computer program also has the capability of computing rocket performance parameters, incident and reflected shock parameters, and Chapman-Jouguet detonation properties.

The method used for the present set of equilibrium composition calculations is based on the concept that reacting chemical systems are in equilibrium when the free-energy of the system is at an extremum value. The specified thermodynamic variables in this report are always pressure and temperature. The appropriate thermodynamic function for these variables is the Gibbs free energy function. As shown by J. B. Callen (1960), this function will possess a minimum. The minimization of the Gibbs free energy is constrained by additional physical conditions. For chemical reaction systems, one appropriate condition is that the mass involved in the reactions be conserved. When a given function possesses an extremum value, which is to be determined under one or more restraining conditions, the most suitable technique for determining this is the method of Lagrangian multipliers.

The alternative method of calculating chemical equilibrium composition is by using sets of simultaneous equilibrium equations in terms of equilibrium constants. The differences between the two methods, with their advantages and disadvantages, has been discussed by F. J. Zeleznik and S. Gordon (1960, 1966, 1968), and S. Gordon (1970). The investivations and results reported in the above-cited papers formed the basis for the computational method and program outlined by S. Gordon and B. J. McBride (1971).

Gordon (1960, 1970) has shown that if 'generalized' methods of solution are used to solve equilibrium problems expressed in either free-energy terms, or by equilibrium constant equations, the resulting sets of equations will reduce to the same number of equations to be solved. However, the 'constant' method has disadvantages not present in the free-energy method. The first of these, as noted by Gordon (1970), is "more bookkeeping". In using the 'constant' approach, the components (amounts) of the reactants and products must be selected, and the possible reactions written in terms of these components. Furthermore, equilibrium constants must be computed. It may be noted that such constants are not always available in the literature for high pressures and/or temperatures. The free-energy method does not require any part of the above-mentioned procedures.

If a chemical species is present in only trace amounts, i.e. "small components," then problems will occur in numerically solving the equilibrium constant equations. In the free-energy method, numerical difficulties are easily circumvented.

When one or more species is present in the form of a condensed phase (liquid and/or solid), the composition problem is very difficult to solve using the 'constant' approach, but is much easier to solve by the free-energy technique.

If non-ideal equations of state are to be used, then the 'constant' method will be very difficult to use, as the generalized solution methods are not very amenable to such modifications.

Partition functions may also be used in solving equilibrium problems. As shown by Lee, Sears, and Turcotte (1963), the equilibrium constant for a given chemical reaction may be expressed in terms of the partition functions of the reacting species. As an example, consider the reaction

$$(1) X + Y = XY.$$

The equilibrium constant for the reaction is

(2)
$$K = \frac{Z_{XY}}{Z_{X}Z_{Y}}$$

where Z_X = the partition function for specie X

 Z_Y = the partition function for specie Y

 Z_{XY} = the partition function for specie XY.

The Gibbs free energy function can also be written in terms of the partition functions:

(3)
$$G_{f} = -k! T \ln \left(\left[\frac{Z_{X}}{n_{X}^{Q}} \right]^{n_{X}^{Q}} \left[\frac{Z_{Y}}{n_{Y}^{Q}} \right]^{n_{Y}^{Q}} \left[\frac{Z_{XY}}{n_{XY}^{Q}} \right]^{n_{XY}^{Q}} \right)$$

where $k^{\,i}\!=\!$ the Boltzmann constant, and $n_X^{\,O},\;n_Y^{\,O},\;$ and $n_{XY}^{\,O}$ are the number of molecules of species X, Y, and XY at equilibrium.

Using the partition functions in the calculation of equilibrium chemical compositions involves computing equilibrium constants, a step which is not necessary in the free-energy method.

In the N.A.S.A. computer program, it is assumed that all gaseous chemical species will behave as ideal gases. For the high temperatures and pressures considered in the present work, the ideal gas law will be a fair approximation to the real behavior of the reacting system. The equation of state for ideal gases is written

(4)
$$\frac{P}{-} = NRT$$
 where
$$N = \sum_{j=1}^{m} n_{j}$$

$$n_{j} = \text{the number of kg-moles of the } j^{\text{th}} \text{ species.}$$

$$m = \text{the number of species in the gas mixture}$$

The following discussion is based in part on the material in the reports of Zeleznik and Gordon (1968) and Gordon and McBride (1971). The Gibbs free-energy function may be defined as

n - m = the number of condensed species.

(5)
$$G_{f} = G_{f}(T,P,n_{j})$$

or equivalently,

$$G_f = U - TS + PV$$

and as

(7)
$$dU = TdS - PdV + \sum_{j}^{m} \mu_{j} dn_{j}$$

then

(8)
$$dG_{f} = -SdT + VdP + \sum_{j} \mu_{j} dn_{j}$$

where the chemical potential (per kg-mole) is defined as

(9)
$$\mu_{j} = \frac{\partial G_{f}}{\partial n_{j}} |_{T,P, n_{k} \neq j}.$$

The Gibbs function is homogeneous in n_j , i.e. the Gibbs function has the property that

(10)
$$G_{f}(T,P,\lambda n_{j}) = \lambda G_{f}(T,P,n_{i}).$$

As previously mentioned, the Gibbs free energy function is appropriate in analyzing systems where the pressure and temperature are of constant value. It was also stated that the function will have a minimum value for a system in equilibrium. A system is considered to be in equilibrium if it is not changing, either mechanically, thermally, or chemically.

To show that the Gibbs function of a system is of minimum value when the system is in equilibrium, consider a system in contact with two 'reservoirs', one at constant temperature, the other at constant pressure. The general conditions for equilibrium are that

$$dU = 0$$

(12)
$$d^2U = 0$$
.

For an equilibrium situation, the total energy of the given system and the two 'reservoirs' must be of minimum value. Hence,

(13)
$$d(U + U_r) = 0$$

where r refers to 'reservoir' conditions. For the constant temperature reservoir,

$$dU_{r,T} = -T_{r}dS$$

and for the constant pressure reservoir,

$$dU_{r,P} = P_r dV.$$

Therefore,

(16)
$$d(U + U_r) = dU - T_r dS + P_r dV$$

(17)
$$d(U + U_r) = d(U - T_r S + P_r V).$$

Since the given system will have a temperature equal to $\mathbf{T}_{\mathbf{r}}\text{,}$ and a pressure equal to $\mathbf{P}_{\mathbf{r}}\text{,}$

(18)
$$d(U + U_r) = d(U - TS + PV)$$

$$d(U + U_r) = dG_f$$

and the condition for equilibrium becomes, for constant temperatures and pressures,

$$dG_{f} = 0.$$

For chemically reacting systems, the equivalent condition is determined by using the Euler relation (Callen, 1960) and substituting into equation (18).

(21)
$$U = TS - PV + \sum_{j=1}^{m} \mu_{j} n_{j}$$

The equilibrium condition then becomes

(22)
$$d(\sum_{j=1}^{m} \mu_{j} n_{j}) = 0.$$

To show that the condition is for a minmum value of the Gibbs function, consider the second derivative of the system energy.

(23)
$$d^{2}(U + U_{r}) = d^{2}U + T_{r}d^{2}S + P_{r}d^{2}V = d^{2}G_{f}$$

and for the reservoirs,

(24)
$$T_r d^2 S = 0$$

(24)
$$T_r d^2 S = 0$$

(25) $P_r d^2 V = 0$

and from equation (12),

(26)
$$d^{2}(U + U_{r}) = d^{2}G_{f} > 0.$$

The Gibbs function thus has a minimum value for a system in equilibrium.

The free-energy per kilogram of gas mixture is, for constant temperature and pressure,

$$g = \sum_{j=1}^{n} \mu_{j} n_{j} .$$

The condition for the mass balance of the reaction system needs to be con-The expression for the mass balance is

(28)
$$\sum_{i=1}^{n} a_{ij} b_{j} - b_{i}^{0} = 0 \quad i = 1, \dots, \ell$$

 a_{ij} = the stoichiometric coefficients (the kg-atoms of element i per kg-mole of specie j) where

 b_i^O = the specified number of kg-atoms of element i per kg of total reactants.

Define a function G as

(29)
$$G = g + \sum_{i=1}^{k} \lambda_{i} (b_{i} - b_{i}^{O})$$

 $b_{i} = \sum_{i=1}^{n} a_{ij} n_{j}$ where

 λ_i = a Lagrangian multiplier

l = the number of chemical elements •

Equation (29) represents the Gibbs Function, which is to be minimized, with the applied constraint. The condition for equilibrium becomes

(30)
$$\delta G = 0 = \sum_{j=1}^{n} \left[\mu_{j} + \sum_{i=1}^{\ell} \lambda_{i} a_{ij} \right] \delta n_{j} + \sum_{i=1}^{\ell} (b_{i} - b_{i}^{O}) \delta \lambda_{i}.$$

If the variations in n, and λ_i are independent, i.e. independent δn_i and $\delta \lambda_i$, then the resulting equilibrium equations are

$$\mu_{j} + \sum_{i=1}^{k} \lambda_{i} a_{ij} = 0$$

$$b_{i} - b_{i}^{O} = 0.$$

It should be noted that λ contains the condition for phase equilibrium, and that the chemical potential of a species is the same for all phases of the species at a given temperature and pressure.

For the ideal gas equation of state, the chemical potential of the j^{th} species may be expressed as

(33)
$$\mu_{j} = \mu_{j}^{O} + RT \ln \left[\frac{n_{j}}{n} \right] + RT \ln \left(P_{a+m} \right)$$

$$j = 1, \dots, m$$

$$j = m + 1, \dots, n$$

where μ_j^O is the chemical potential in the standard state, and P_{atm} is pressure given in terms of atmospheres.

The equations to be solved by Newton-Raphson iteration for the n_j , λ_i , and the numbers of moles, are listed in the report by Gordon and McBride. Details on initial estimates for the chemical composition, criteria for the covergence of the numerical solutions, tests for the possibility of condensed phases, and phase transitions, and mathematical singularities are also to be found in the report.

The thermodynamic data provided in the report with the program listing are mostly from the JANAF thermochemical property tables. For each chemical species, the specific heat for constant pressures, the enthalpy, and the entropy were computed from the following equations.

(35)
$$\frac{c_{P}^{o}}{R} = a_{1} + a_{2}T + a_{3}T^{2} + a_{4}T^{3} + a_{5}T^{4}$$

(36)
$$\frac{H_{T}^{0}}{RT} = a_{1} + \frac{a_{2}}{2} + \frac{a_{3}}{3} + \frac{a_{4}}{4} + \frac{a_{5}}{5} + \frac{a_{6}}{5}$$

(37)
$$\frac{s_{T}^{o}}{R} = a_{1} \ln T + a_{2} T + \frac{a_{3}}{2} T^{2} + \frac{a_{4}}{3} T^{3} + \frac{a_{5}}{4} T^{4} + a_{7}$$

where the a are 'least-squares' coefficients, which are provided in the program data deck. Two sets of coefficients are listed, for the temperature ranges 300 $^{\circ}$ K to 1000 $^{\circ}$ K, and from 1000 $^{\circ}$ K to 5000 $^{\circ}$ K. The use of either set of coefficients will give identical results at a temperature of 1000 $^{\circ}$ K.

A comparison may be **STHUBBARROW MOISCUDSIO** of the property tion and the results of A. S. Kesten (1917). A sammary of the reality compared is given in the following tools.

The results of the present set of computations for the decomposition of hydrazine are presented in both graphical and tabular form. The temperature range, from 300 °K to 1000°K, is divided into increments of 25 °K. The pressure range, from 1.0 atm to 610 atm, is divided into increments of 10 atm. The tabulated results are given in the appendix. The graphical presentation consists of three sets of figures, each set consisting of three figures, one for each of the reaction product species. The first set of figures are plots of mole-fraction (of H2, N2, or NH3) vs temperature with pressure as the parameter. The second set of figures consists of graphs of mole-fraction vs pressure, with temperature as the parameter. The last set of figures consists of plots of temperature vs pressure with mole-fractions as the parameter.

For low temperatures and pressures, the decomposition products are approximately 2% H₂, 20% N₂, and 78% NH₃. For pressures greater than 50 atm, no H₂ is produced, and the product mixture is 20% N₂ and 80% NH₃. For higher temperatures (near 1000 °K) and low pressures, no NH₃ is left after the decomposition reaction has finished, and the product mixture consists of approximately 33% N₂ and 67% H₂. For high temperatures and pressures, the product mixture consists of approximately 59% H₂, 32% N₂, and 9% NH₃.

Trace amounts of NH and NH $_2$ may be found in the gaseous product mixture for elevated pressures and temperatures, but the amounts of these species are negligible.

It is clear that the effect of increasing the pressure, for a given temperature, is to decrease the amount of $\rm H_2$ (and subsequently the amount of $\rm N_2$) in the product mixture, and increase the amount of $\rm NH_3$.

At low temperatures, the over-all decomposition reaction may be special represented by

(i) the contract of the property of the contract of the contr

A comparison may be made between the results of the present computation and the results of A. S. Keston (1967). A summary of the results compared is given in the following table.

			computations.

P, atm	T , ° K	Mole Fr	action (Keston) NH ₃	Mole Fra H ₂	ction (pres	sent work).
7.59	1000	.55 (.59)	.31	.14	.665	.335	.0012
32.6	1054	.55 (.56)	.31	.14	.660	.330	.0059
71.0	1054	.57 (.54)	.32	.12 (.15)	.665	.328	.0110

The numbers in () indicate experimentally-determined values (Keston, 1967).

A comparison between the sets of results shows good agreement as to the amount of nitrogen produced. The discrepancies between the results are for the mole-fractions of hydrogen and ammonia. For both the theoretical predictions and experimental findings, Keston's results indicate that fewer moles of hydrogen will be produced, and many more moles of ammonia. Qualitatively, both investigations showed that an increase in the pressure for a given temperature will result in the production of more ammonia. It should be noted, however, that Keston's results are for steady flow of hydrazine across a Shell 405 catalyst bed. The data he presents do not appear to have reached zero rate of change with axial position in the catalyst bed and thus can be expected to show the lower $\rm H_2$ concentrations than would equilibrium calculations; this is because most of the $\rm H_2$ is produced by an endothermic decomposition of NH3 which proceeds at a finite rate, evidently, slower than the transit time of the gas through the catalyst bed.

Two factors, in addition to the basic thermodynamic data used in the computation, certainly affect the accuracy of these results: The first is the condensation of one or more of the species at certain conditions and the second is the real gas behavior of the various species. One would expect that for the temperatures and pressures considered in the present investigation $\rm H_2$ and $\rm N_2$ behave substantially as perfect gases, far removed from saturation. $\rm NH_3$, on the other hand, may easily saturate (see Table 6). Thus one should also consider the non-ideal nature of $\rm NH_3$ for pressures and temperatures immediately above saturation in estimating the mole fraction of the various species present.

Table 6. Saturation curve of ammonia.

T∿°K	P∿atm
195.42 (T.P.) 200 210 220 230 240 250 260 270 280 290 300 310 320 330 340 350 360 370 380 390 400 405.6 (C.P.)	0.05997 0.08553 0.1756 0.3347 0.5982 1.0124 1.6320 2.526 3.769 5.446 7.652 10.485 14.063 18.49 23.90 30.42 38.19 47.40 58.14 70.60 84.94 101.5

Figure 10 compares the ideal NH_3 gas pressure for given density and temperature with the real gas pressure (Din, 1956) at the same conditions. The trends shown in the figure can be explained as follows: As the pressure increases at constant temperature the energy in the NH₃ molecule first tends to migrate from the translational degrees of freedom to the rotational or vibrational degree of freedom. Consequently the molecule transfers less momentum (exerts a smaller pressure) as the result of elastic collisions with other molecules than would be expected from a knowledge of its energy content. As one attempts to crowd more and more molecules of the same energy content into a given volume, collisions finally become so numerous that all the atoms in the molecule become involved more or less simultaneously, giving one the appearance of reducing the activity in the vibrational degree of freedom. This process is most apparent where the total energy content is small, e.g. at low temperatures, just above condensation. In the worst case the real gas pressure is about a factor of three lower than the ideal gas pressure for the same specific volume and temperature.

The results cannot be applied exactly fo a mixture. The total pressure gives an indication of the crowding of the molecules but N_2 and H_2 molecules may induce a somewhat different response from NH $_3$ than NH $_3$ molecules on NH $_3$. Generally, one would expect that if the NH $_3$ in a mixture would suddenly go from a perfect gas to a real gas, the total pressure would drop somewhat (in proportion to the change in NH $_3$ pressure for the prevailing temperature and specific volume) and that the volume fractions of N_2 and H_2 would thereby increase somewhat. However, the total pressure is then somewhat lower and the NH $_3$ mole-fraction would then tend to increase. The net result is probably little different from that given by the perfect gas computations.

In comparing equilibrium computations with results for catalytic beds it is instructive to recall that an adiabatic catalyst can promote the redistribution of the energy brought to a reaction by the constituents but it does not add or withdraw energy (unless heated or cooled). In the case of hydrazine the initial reaction

(3)
$$3N_2H_4 \rightarrow 4NH_3 + N_2$$

3.300.04

is exothermic. If the energy released in this process is all used to effect the endothermic decomposition of NH $_3$ to N $_2$ and H $_2$ then when the temperatures and pressures of all the reaction products equals that of the incoming N $_2$ H $_4$, the mole fractions of the product species will be as given by the equilibrium calculation. Viewed in this light the equilibrium calculations provide an upper bound for the effectiveness of a catalytic reactor.

In the real world, of course, things are not so simple. The reactor may have to be cooled, reducing the amount of heat available to decompose the NH3. The reactions may not be complete when the constituents leave the reactive area and enter a region of lower pressure and lower temperature. In this region the reactions may freeze leaving the mole-fractions what they were leaving the reactive area or they may go to completion at the lower temperatures and pressures. In the latter case the NH3 concentrations will be greater and the H2 concentration lower than if the reactions froze. Thus to get large quantities of H2 it is necessary to keep the reactants not long enough for the NH3 to decompose completely and then cool the mixture rapidly. Another factor to be considered in designing a reactor for H2 production from N2H4 is the data of Sayer which indicated that a rhodium catalyst is more effective at promoting direct H2 formation than the iridium (Shell 405) catalyst which seems to favor ammonia production as an intermediate step.

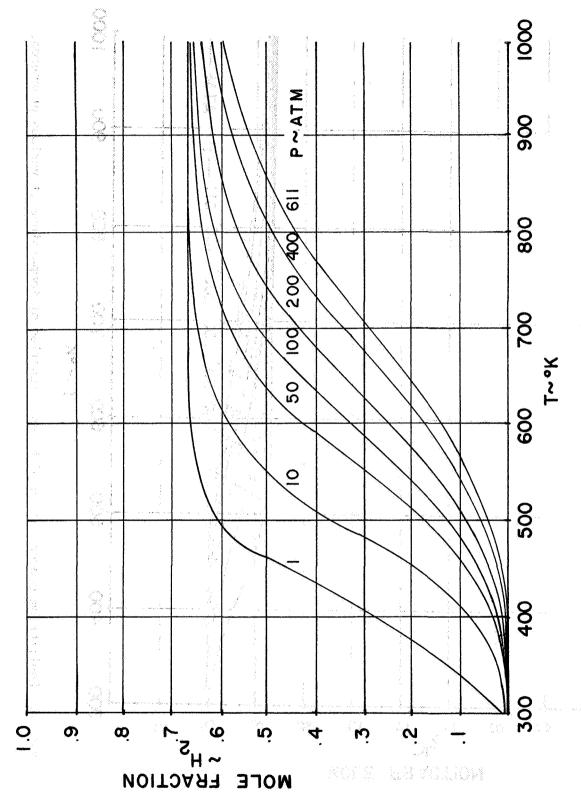


Figure 1. Equilibrium concentration of H₂ as a function of temperature at various pressures.

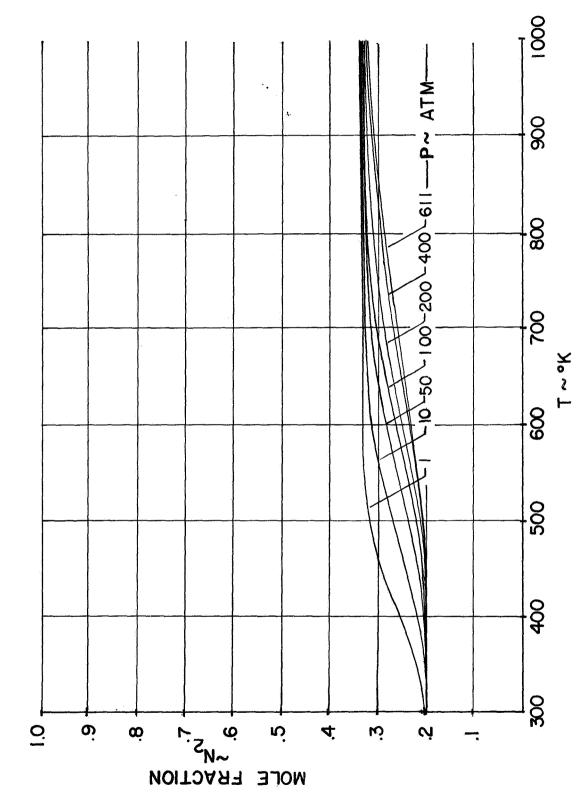


Figure 2. Equilibrium concentration of N_2 as a function of temperature at various pressures.

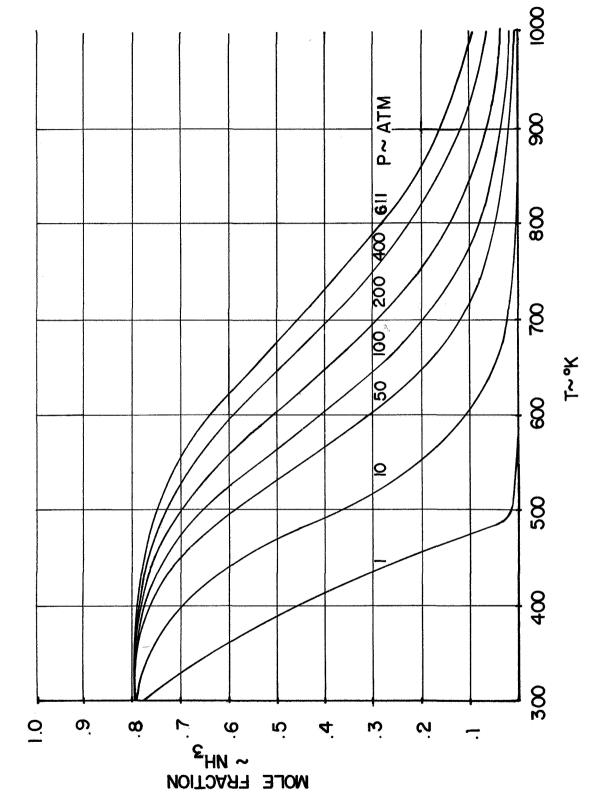
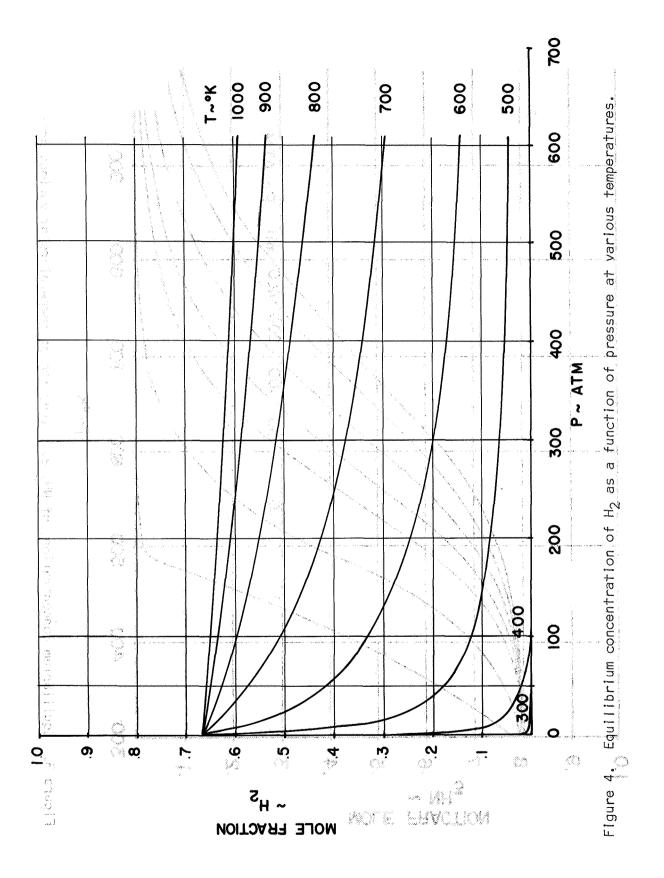


Figure 3. Equilibrium concentration of ${
m NH}_{
m 3}$ as a function of temperature at various pressures.



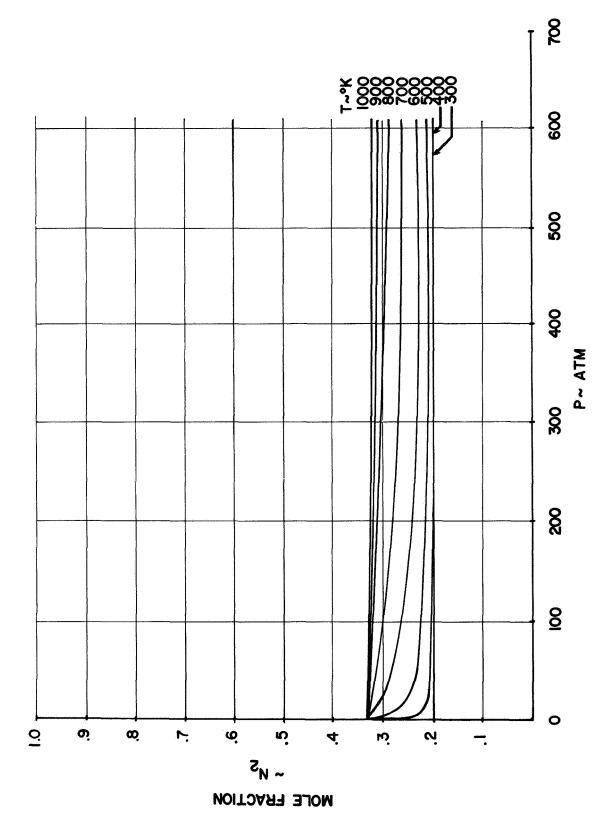


Figure 5. Equilibrium concentration of ${\sf N}_2$ as a function of pressure at various temperatures.

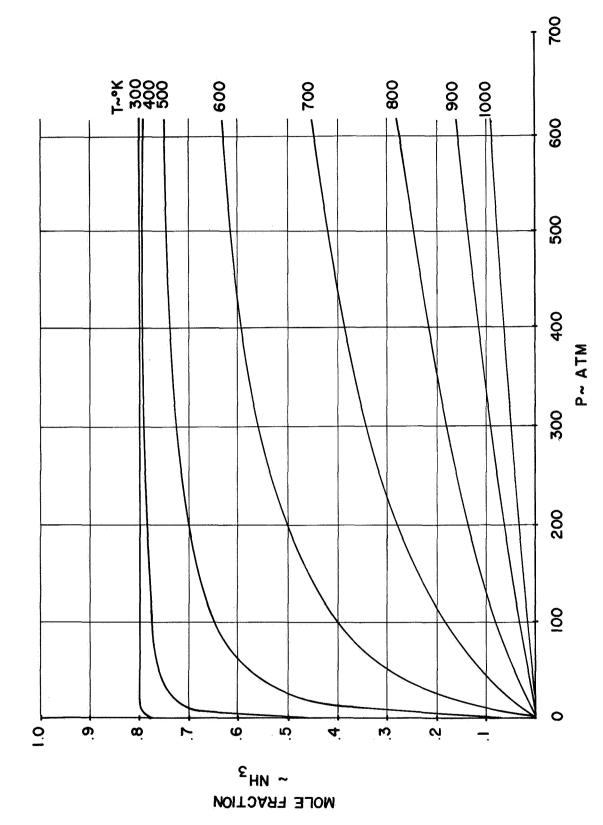


Figure 6. Equilibrium concentration of NH_{γ} as a function of pressure at various temperatures.

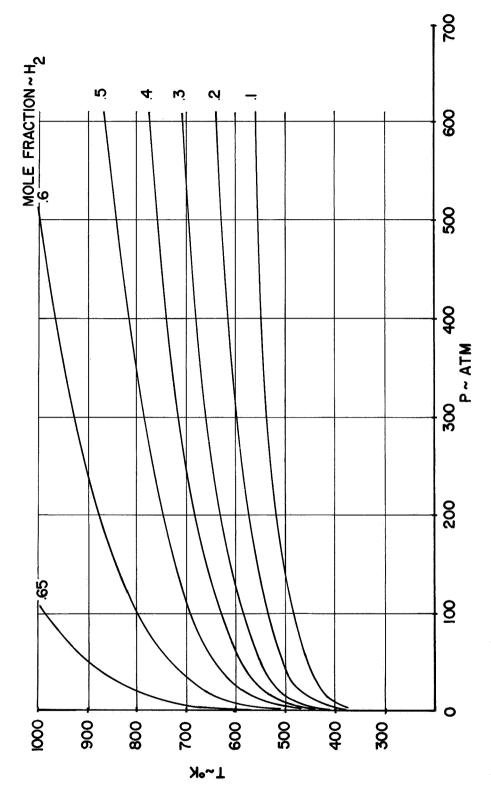


Figure 7. Variation of H_2 pressure with temperature for constant mole fractions.

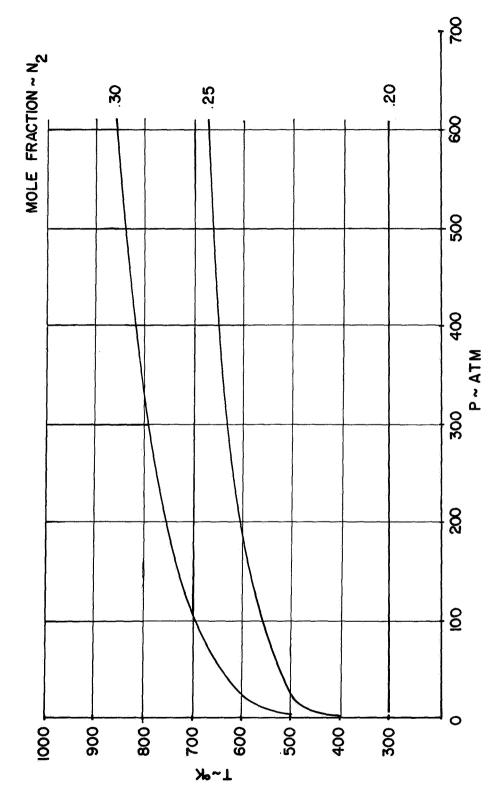


Figure 8. Variation of $N_{
m 2}$ pressure with temperature for constant mole fractions.

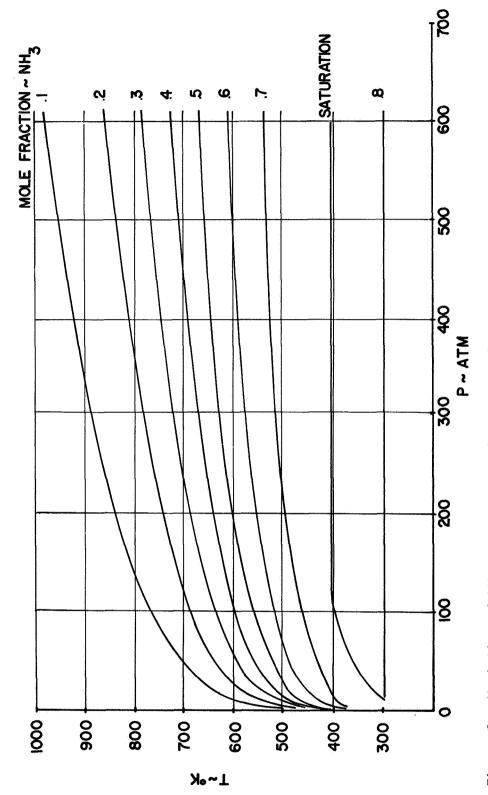


Figure 9. Variation of NH_3 pressure with temperature for constant mole fractions.

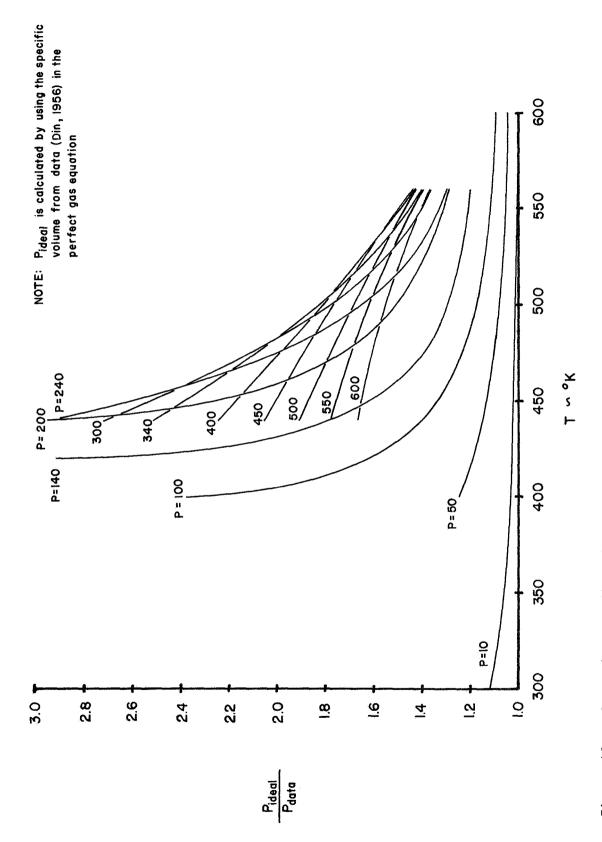


Figure 10a. Comparison of the behavior of NH_3 with a perfect gas.

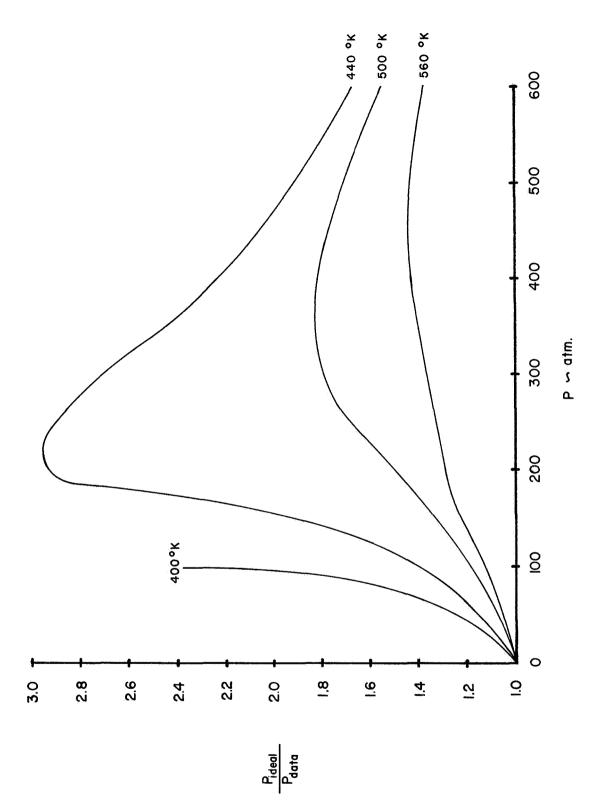


Figure 10b. Comparison of the behavior of NH3 with a perfect gas.

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APPENDIX A - Nomenclature

- A Helmholtz free energy
- $a_{i\,i}$ stoichiometric coefficient
- $b_i = \sum_{i} a_{ij}^n j$
- b^o specified number of kg-atoms of element i per kg of total reactants
- c_{v} constant volume specific heat
- ${\bf c}_{\bf p}$ constant pressure specific heat
- ${\rm G_f}$ Gibbs free energy function
- g Gibbs free energy per kg of gas mixture
- H enthalpy
- K equilibrium constant, general reaction equation
- k reaction rate constant
- k' Boltzmann constant
- k, reaction rate constant for first order reaction
- ${\bf k}_2$ reaction rate constant for second order reaction
- m number of species in the gas mixture
- N total number of moles
- n number of species in the reacting system
- n_X number of moles of specie X
- n_{Υ} number of moles of specie Y
- n_{XY} number of moles of reaction product XY
- P pressure
- R gas constant
- S entropy

- T temperature
- t time
- U internal energy
- V specific volume
- Z partition function
- Z_{χ} partition function for specie X
- Z_{V} partition function for specie Y
- Z_{XY} partition function for the reaction product XY
- α thermal expansion coefficient
- β compressibility
- λ Lagrangian multiplier
- u chemical potential
- ρ density

Superscripts

o quantity at equilibrium

Subscripts

- c value of quantity at the critical point
- i chemical specie indexing number
- j chemical specie indexing number
- r 'reservoir' condition
- T value for specified temperature
- [0₂] brackets containing chemical symbols refer to concentrations of those chemical species.

Note: Duplicate symbols and/or other symbols appearing in the report are defined in the section where used.

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6596 0.0085 0.3319 625 0.6054 0.0735 0.3211 625 0.5595 0.1286 0.311 6618 0.0058 0.3324 650 0.6229 0.0525 0.3246 650 0.6229 0.0525 0.3270 675 0.679 0.0705 0.031 6642 0.0029 0.3326 675 0.6435 0.0278 0.3287 700 0.6229 0.0705 0.3244 0.0705 0.3279 6642 0.0021 0.3326 725 0.6495 0.0206 0.3239 725 0.6229 0.0526 0.3249 6649 0.0021 0.3331 775 0.6495 0.0119 0.3313 775 0.6413 0.0329 775 0.6413 0.0324 0.3244 0.0324 0.3246 6659 0.0012 0.3331 775 0.6413 0.0119 0.3318 800 0.6514 0.0114 0.3324 800 0.6516 0.0005 0.3324 800 0.6516 <th< td=""><td>. 656</td><td>10.</td><td>.331</td><td>0</td><td>.580</td><td>.103</td><td>.316</td><td>0</td><td>. 52</td><td>.173</td><td>.304</td></th<>	. 656	10.	.331	0	.580	.103	.316	0	. 52	.173	.304
6618 0.0058 0.3224 650 0.6229 0.3276 0.3270 657 0.0079 0.00951 0.317 6633 0.0041 0.3326 675 0.6351 0.0379 0.3270 675 0.6079 0.0705 0.324 6642 0.0029 0.3328 700 0.6435 0.0278 0.3287 700 0.6229 0.0526 0.3243 6649 0.0021 0.3330 725 0.6495 0.0206 0.3299 725 0.6537 0.0366 6659 0.0012 0.3331 775 0.6537 0.0119 0.3313 775 0.6473 0.0322 0.328 6659 0.0012 0.3332 800 0.6567 0.0119 0.3321 775 0.6473 0.0322 0.328 6660 0.0007 0.3332 825 0.6608 0.0003 0.3324 850 0.6518 0.0004 0.3324 850 0.6518 0.0004 0.3324 850 0.6518 0.0004	. 659	00.	.331	2	. 605	.073	.321	2	.55	.128	.311
6633 0.0041 0.3326 675 0.6351 0.0379 0.3270 675 0.6079 0.0705 0.324 6642 0.0029 0.3328 700 0.6435 0.0278 0.3287 700 0.6229 0.0526 0.324 6649 0.0021 0.3331 725 0.6435 0.0206 0.3299 725 0.6337 0.0326 0.326 6659 0.0012 0.3331 775 0.6415 0.0331 775 0.6415 0.0331 0.328 6659 0.0012 0.3332 800 0.6567 0.0119 0.3313 775 0.6473 0.032 0.328 6660 0.0007 0.3332 825 0.6506 0.0003 0.3324 825 0.6676 0.0181 0.3324 825 0.6676 0.0014 0.3324 825 0.6676 0.0014 0.3324 825 0.6674 0.0143 0.3324 6663 0.0004 0.3333 0.6628 0.0046 0.3324	. 661	9	.332	5	.622	.052	.324	5	• 58	• 095	.317
.6642 0.0029 0.3328 700 0.6435 0.0278 0.3287 700 0.6229 0.0526 0.3264 .6649 0.0021 0.3330 725 0.6495 0.0206 0.3299 725 0.6337 0.0396 0.326 .6659 0.0012 0.3331 775 0.6567 0.0119 0.3313 775 0.6415 0.0332 0.328 .6659 0.0012 0.3332 800 0.6567 0.0119 0.3313 775 0.6473 0.0232 0.328 .6659 0.0009 0.3332 800 0.6571 0.0181 0.329 .6650 0.0007 0.3332 825 0.6618 0.0073 0.3321 825 0.6516 0.0143 0.332 .6664 0.0005 0.3332 825 0.6618 0.0058 0.3324 850 0.6517 0.0114 0.331 .6664 0.0004 0.3333 0.6645 0.0038 0.3324 850 0.6510 0.	. 663	90.	. 332	-	.635	.037	.327		99.	.070	.321
.6649 0.0021 0.3330 725 0.6495 0.0206 0.3299 725 0.6337 0.0396 0.328 .6653 0.0016 0.3331 750 0.6415 0.0301 0.328 .6656 0.0012 0.3331 775 0.6567 0.0119 0.3313 775 0.6473 0.0232 0.328 .6659 0.00012 0.3332 800 0.6590 0.0092 0.3318 800 0.6516 0.0181 0.328 .6660 0.0007 0.3332 825 0.6618 0.0073 0.3321 825 0.6516 0.0181 0.332 .6663 0.0006 0.3332 875 0.6618 0.0046 0.3324 850 0.6571 0.0143 0.331 .6663 0.0006 0.3333 975 0.66571 0.0014 0.0092 0.3324 875 0.6679 0.0092 0.3324 875 0.6679 0.0092 0.3324 975 0.6679 0.0092 0.3324 <	.664	00.	.332	0	.643	.027	.328	0	•62	.052	.324
.6653 0.0016 0.3331 750 0.6557 0.0156 0.3313 775 0.6415 0.0301 0.328 .6656 0.0012 0.3331 775 0.6567 0.0119 0.3313 775 0.6473 0.0232 0.329 .6659 0.0012 0.3332 800 0.6590 0.0092 0.3321 800 0.6516 0.0181 0.332 .6660 0.0007 0.3332 825 0.6606 0.0073 0.3321 825 0.6547 0.0181 0.330 .6662 0.0006 0.3332 850 0.6618 0.0074 0.3324 850 0.6571 0.0143 0.332 .6663 0.0006 0.3333 975 0.6628 0.0046 0.3324 850 0.6571 0.0114 0.332 .6663 0.0004 0.3333 925 0.6638 0.3324 875 0.6574 0.0114 0.332 .6664 0.0003 0.3333 925 0.6641 0.0026	- 664	00	333	N	649	.020	.329	2	• 63	• 039	.326
.6656 0.0012 0.3331 775 0.6567 0.0119 0.3313 775 0.6567 0.0134 0.3321 800 0.6516 0.0181 0.330 .6659 0.0009 0.3321 825 0.6516 0.0181 0.330 .6662 0.0007 0.3322 825 0.6618 0.0073 0.3324 850 0.6571 0.0143 0.330 .6662 0.0006 0.3332 875 0.6628 0.0058 0.3324 875 0.6571 0.0144 0.331 .6663 0.0006 0.3333 875 0.6628 0.0046 0.3326 875 0.6571 0.0114 0.331 .6663 0.0004 0.3333 900 0.6635 0.0038 0.3327 900 0.6504 0.0075 0.332 .6664 0.0003 0.3333 925 0.6645 0.0022 0.3332 925 0.6645 0.0022 0.3332 975 0.6649 0.0022 0.3333 1000 0	. 665	00	.333	S	.653	.015	.330	5	•64	.030	.328
6659 0.0009 0.3318 800 0.6516 0.0181 0.330 6660 0.0007 0.3321 825 0.6606 0.0073 0.3321 825 0.6547 0.0143 0.330 6660 0.0007 0.058 0.3324 850 0.6571 0.0144 0.331 6663 0.0006 0.3333 875 0.6628 0.0046 0.3326 875 0.6536 0.0114 0.331 6663 0.0006 0.3333 970 0.6638 0.0032 0.3328 975 0.6649 0.0032 0.3329 975 0.6649 0.0022 0.3329 975 0.6649 0.0022 0.3339 975 0.6649 0.0022 0.3330 975 0.6631 0.0036 0.3330 975 0.6631 0.0036 0.3329 975 0.6636 0.0043 0.332 6664 0.0002 0.3330 975 0.6649 0.0022 0.3330 975 0.6631 0.0036 0.332 <	-665	00	.333	-	.656	.011	.331	-	• 64	.023	.329
6660 0.0007 0.3321 825 0.6547 0.0143 0.330 6662 0.0006 0.3332 850 0.6618 0.0058 0.3324 850 0.6571 0.0114 0.331 6663 0.0006 0.3333 875 0.6628 0.0046 0.3326 875 0.6590 0.0114 0.331 6663 0.0004 0.3333 900 0.6635 0.0038 0.3327 900 0.6590 0.0092 0.332 6664 0.0003 0.3333 925 0.6641 0.0031 0.3328 925 0.6615 0.0062 0.332 6664 0.0003 0.3333 950 0.6645 0.0026 0.3329 950 0.6615 0.0062 0.3339 6664 0.0002 0.3333 975 0.6649 0.0022 0.3330 975 0.6634 0.0036 0.0649 0.03330 1000 0.6636 0.0036 0.0332 6665 0.0002 0.3333 1000 <td>. 665</td> <td>00</td> <td>.333</td> <td>0</td> <td>.659</td> <td>•000</td> <td>.331</td> <td>0</td> <td>• 65</td> <td>.018</td> <td>.330</td>	. 665	00	.333	0	.659	•000	.331	0	• 65	.018	.330
.6662 0.0006 0.3332 850 0.6618 0.0056 0.3324 850 0.6571 0.0114 0.331 .6663 0.0005 0.3333 875 0.6636 0.0046 0.3326 875 0.6590 0.00092 0.331 .6663 0.0004 0.3333 900 0.6635 0.0038 0.3327 900 0.6604 0.0075 0.332 .6664 0.0003 0.3333 925 0.6641 0.0031 0.3328 925 0.6615 0.0062 0.332 .6664 0.0003 0.3333 950 0.6645 0.0026 0.3329 950 0.6615 0.0062 0.332 .6665 0.0002 0.3333 1000 0.6649 0.0022 0.3330 975 0.6634 0.0036 0.0332 .6665 0.0002 0.3333 1000 0.6651 0.0033 0.0364 0.0036 0.332	.666	9	.333	2	.660	.007	.332	2	• 65	.014	.330
.6663 0.0005 0.3333 875 0.6628 0.0046 0.3326 875 0.6590 0.0092 0.331 .6663 0.0004 0.3333 900 0.6635 0.0038 0.3327 900 0.6604 0.0075 0.332 .6664 0.0003 0.3333 925 0.6641 0.0031 0.3328 925 0.6615 0.0062 0.332 .6664 0.0003 0.3333 975 0.6645 0.0026 0.3339 975 0.6649 0.0022 0.3330 975 0.6649 0.0018 0.3330 975 0.6651 0.0018 0.3330 1000 0.6636 0.0036 0.332 .6665 0.0002 0.3333 1000 0.6651 0.0018 0.3330 1000 0.6636 0.0036 0.332	666	00.	.333	S	.661	.005	.332	S	• 65	.011	.331
.6663 0.0004 0.3333 900 0.6635 0.0038 0.3328 900 0.6604 0.0075 0.332 .6664 0.0003 0.3333 925 0.6641 0.0031 0.3328 925 0.6615 0.0062 0.332 .6664 0.0003 0.3333 950 0.6645 0.0026 0.3329 950 0.6624 0.0051 0.332 .6665 0.0002 0.3333 975 0.6649 0.0022 0.3330 975 0.6651 0.0018 0.3330 1000 0.6634 0.0036 0.333 .6665 0.0002 0.3333 1000 0.6651 0.0018 0.3330 1000 0.6636 0.0036 0.332	. 666	0	.333	-	.662	.004	.332	-	• 65	600	.331
.6664 0.0003 0.3333 925 0.6641 0.0031 0.3328 925 0.6615 0.0062 0.332 .6664 0.0003 0.3333 950 0.6645 0.0026 0.3329 950 0.6624 0.0051 0.332 .6665 0.0002 0.3333 975 0.6649 0.0022 0.3330 975 0.6631 0.0043 0.332 .6665 0.0002 0.3333 1000 0.6651 0.0018 0.3333 1000 0.6636 0.0036 0.332	. 666	00	.333	0	.663	.003	.332	0	• 66	.007	.332
.6664 0.0003 0.3333 950 0.6645 0.0026 0.3329 950 0.6624 0.0051 0.332 .6665 0.0002 0.3333 975 0.6649 0.0022 0.3330 975 0.6631 0.0043 0.332 .6665 0.0002 0.3333 1000 0.6651 0.0018 0.3330 1000 0.6636 0.0036 0.332	.666	00	.333	N	\$999	.003	.332	2	•66	.006	.332
.6665 0.0002 0.3333 975 0.6649 0.0022 0.3330 975 0.6631 0.0043 0.332 .6665 0.0002 0.3333 1000 0.6651 0.0018 0.3330 1000 0.6636 0.0036 0.332	. 666	00.	.333	5	+999*	.002	.332	5	• 66	.005	.332
.6665 0.0002 0.3333 1000 0.6651 0.0018 0.3330 1000 0.6636 0.0036 0.332	. 666	90.	.333	-	.664	.002	.333	-	• 66	.004	.332
	.666	.00	.333	9	• 665	.001	.333	0	• 66	• 003	32

PRESSURE = 50.0 ATM.	P H2 NH3	00 0.0014 0.7983 0.200	25 0.0037 0.7956 0.200	50 0.0083 0.7900 0.201	75 0.0170 0.7796 0.203	00 0.0315 0.7621 0.206	25 0.0540 0.7352 0.210	50 0.0860 0.6968 0.217	75 0.1281 0.6462 0.225	00 0.1795 0.5845 0.235	25 0.2379 0.5146 0.247	50 0.2995 0.4405 0.259	75 0.3606 0.3672 0.272	00 0.4176 0.2989 0.283	25 0.4679 0.2386 0.293	50 0.5102 0.1877 0.302	75 0.5447 0.1464 0.308	00 0.5719 0.1138 0.314	25 0.5930 0.0884 0.318	0 0.6091 0.0690 0.	75 0.6215 0.0542 0.324	00 0.6309 0.0429 0.326	25 0.6381 0.0342 0.327	50 0.6437 0.0275 0.328	75 0.6480 0.0224 0.329	00 0.6514 0.0183 0.330	25 0.6540 0.0152 0.330	50 0.6561 0.0126 0.331	75 0.6578 0.0106 0.331	
		.200	.200	.201	.203	.207	.212	.219	.229	.240	.252	.265	.277	.289	.298	.306	.312	317	.321		.325	.327	.328	.329	.330	.330	.331	.331	.331	
40.0 ATM.	NH3	.798	,794°	.788	.776	.756	. 725	.682	.626	.558	.484	.407	.332	.265	.208	.161	.123	.095	.073	0.0568	,044	.034	.027	.022	.018	.014	.012	.010	.008	
II	7	.001	.004	600-	.019	.036	.062	.098	.144	.201	.263	.327	.389	445	.493	.532	.563	.587	.605	O	.629	.637	.643	.648	.651	.654	.656	.658	.659	
PRESSURE	3	300	2	5		0	2	S	-	0	2	5	-	0	2	5	-	0	2	S	Sec.	0	2	5		0	N	5	-	
		0.2004	. 20	.20	.20	.20	.21	.22	•23	.24	.25	.27	.28	.29	.30	.31	33	32	32		.32	.32	.32	33	(C)	60	60	.03	E.C.	
30.0 ATM.	工	-	.79	.78	77.	.74	7	99.	.59	52	44	36	28	.22	.17	E	60.	.07	0.0	9	.03	.02	.02	TO.	TO.	0	000	00:	00.	! :
)) ·	75	.002	.005	.011	.023	.043	.073	.115	.168	.231	.297	.364	.425	.479	. 523	.557	. 584	.604	.619	0.6301	.638	.644	•649	.652	.655	.657	.659	.660	.661	
PRESSURE	TEMP	300	325	350	375	400	425	450	475	200	525	550	575	009	625	650	675	700	725	750	775	800	825	850	875	006	925	950	975	

ATM.	Z	7 0.200	8 0.200	6 0.201	0 0.2025	0 0.204	7 0.208	2 0.213	4 0.219	1 0.228	9 0.237	1 0.248	6 0.260	9 0.271	4 0.282	7 0.291	9 0.299	3 0.306	4 0.311	1 0.316	4 0.319	2 0.322	5 0.324	6 0.326	8 0.327	7 0.328	8 0.329	9 0.330	8 0.330	2 0.331
80.0 A	I	0.79	0.79	0.79	0.785	0.77	0.75	0.72	0.68	0.63	0.57	0.50	0.43	0.37	0.30	0.25	0.20	0.16	0.12	0.10	0.08	0.06	0.05	0.04	0.03	0.02	0.02	0.0	0.01	0.01
ESSURE =	H2	00.	8	8	0.0125	.02	•04	•06	60.	.13	.18	.24	.30	35	.41	.45	64.	.53	. 55	.58	.59	.61	.62	• 63	M	•64	•64	•65	5	S
PRES	2	0	S	5	375	0	N	S	-	0	N	5	-	0	2	S	-	0	2	5	-	0	2	S	-	0	2	5	~	0
_	N 2	. 200	. 200	.201	0.2027	. 205	.208	.214	.221	.230	.240	.251	.263	.274	.285	.294	.302	.308	.314	.318	.321	.323	.325	.327	.328	.329	.329	.330	.330	.331
70.0 ATM.	NH3	. 798	. 796	. 792	0.7836	.769	.747	• 715	.672	.619	.557	.488	.418	.350	.287	.231	.185	.146	215	160.	.072	.058	.046	.037	.030	.025	.021	.017	.014	.012
SSURE =		.001	.002	•000	0.0136	.025	.043	.070	.105	.150	.202	.259	.317	.374	.427	.473	.512	. 544	.570	.590	909°	.618	.627	.635	.641	.645	649	.652	.654	.656
PRESS	T	0	O	5	375	0	2	5	-	0	2	S	-	0	S	5	-	0	S	5	-	0	N	5	-	0	N	5	-	1000
•		. 200	.200	.201	0	.205	.209	.215	.223	.232	.243	.255	.267	.278	. 289	.298	.305	.311	.31e	.319	.322	.324	.326	.327	.328	.329	.330	.330	.331	.331
60.0 ATM	NH3	• 79	.79	.79	. 78	.76	.74	.70	99.	.60	533	.46	.39	.32	.26	.21	16	.13	100	.08	• 06	.05	40.	.03	.02	• 02	.01	.01	0	0.0108
SSURE =	2	00.	00.	00.	10.	. 02	40.	.07	-	16	• 2I	-27	.33	.39	44	640	. 52	.55	.58	. 59	.61	.62	.63	. 63	.64	49.	. 65	. 65	•65	0.6577
PRESS	I	30	N	5	-	0	2	5	-	0	N	S	-	0	N	5	-	0	N	5	-	0	N	5	-	Ö	N	S	-	1000

110.0 ATM.	H3 N	.7990 0.20	.7974 0.20	.7940 0.20	.7878 0.20	.7772 0.20	.7605 0.20	360 0.2	.7025 0.21	.6594 0.22	.6074 0.23	.5482 0.24	4845 0.25	.4194 0.26	.3562 0.27	.2975 0.28	.2451 0.29	.1999 0.30	.1620 0.30	.1308 0.31	.1057 0.31	.0855 0.31	.0695 0.32	.0568 0.32	.0467 0.32	.0386 0.32	.0322 0.32	.0270 0.32	.0228 0.32	***
11	H2	*0008	.0022	.0050	01	.0190	.0329	.0533	.0812	.1171	.1605	.2098	•2629	.3171	.3698	.4187	•4624	.5001	5317	.5576	86	.5954	.6087	.6193	.6278	.6345	•6388	.6442	.6477	L
PRESSURE	Z	300	2	5		0	2	S		0	2	5	-	0	N	3	-	0	S	5		\bigcirc	N	3	_	0	S	S	-	•
	N2	-200	.200	.201	.202	.204	.207	0.2113	.217	.224	.233	.243	.254	.265	.276	.286	.294	.302	.308	.313	317	.320	.322	.324	.326	.327	.328	.329	.329	•
100.0 ATM.	I	. 798	197	. 793	.787	.775	.758	(1)	•696	.651	.597	.536	.471	.405	.341	.283	.231	.187	.151	121.	.097	.079	•064	.052	,042	.035	.029	.024	.020	•
a	H2	0000	.002	.005	.010	.020	.034	0	.086	.123	.168	.219	.274	.329	.382	.430	.473	.510	.540	.565	.585	.600	.613	.623	.631	.637	.642	.646	649	4
PRESSURE	I	300	N	5		0	S	5	-	0	N	5	-	0	N	5	-	0	S	5	-	0	2	5	-	0	2	5	-	4
	N2	. 20	.20	. 20	.20	.20	. 20	0.2121	.21	• 22	.23	.24	.25	.26	.27	.28	.29	30	.31	.31	5	32	3	32	32	32	.32	32	33	(
90.0 ATM.	NH3	.798	797	.793	.786	.774	.755	0.7276	069.	.642	.586	. 522	.455	.389	325	.267	.217	175	.140	.112	.089	.072	.058	.047	038	.032	.026	.022	.018	•
11	2	.001	.002	.005	.011	.021	.037	0.0603	.091	.131	.178	.231	.286	.342	.395	. 443	.485	.520	. 549	.573	.591	.606	.618	.627	634	.640	,644	.648	.651	3
PRESSURE	Σ	300	N	S	-	0	N	5	-	0	N	S	-	0	N	S	-	0	N	ເດ	-	0	\sim	S	-	\circ	N	S	-	4

	NZ	• 200	.200	• 200	0.2017	• 203	.205	.209	.214	.220	.228	.237	.247	.257	.267	.277	•286	.294	.301	.307	.312	.316	•319	.321	*323	. 325	.326	-327	• 328	•329
40.0 ATM.	I	. 199	197.	. 794	0.7896	. 780	• 766	• 745	.715	.677	.630	.576	.517	.455	.393	.334	.280	.232	190	-156	.127	. 104	.085	.070	.057	.048	.040	.033	.028	•024
)I	N	0000	.001	. 004	0.0086	.016	.028	.045	.070	.102	.140	.186	.235	.287	.338	.388	.433	.473	.507	.536	. 560	.579	• 595	. 608	.618	.626	.633	.638	.642	•646
PRESSURE	TEMP	300	325	350	375	400	425	450	475	500	525	550	575	909	625	650	675	700	725	750	775	800	825	850	875	900	925	950	975	1000
	N2	.200	.200	.200	0.2018	.203	.205	.209	.214	.221	.229	.238	. 248	.259	.269	.279	.288	.296	.303	.308	.313	.317	.320	.322	.324	.325	.327	.328	.328	•329
30.0 ATM.	NH3	662.	197	* 794	0.7891	.779	. 764	. 742	•711	.672	.623	.568	.507	.444	.382	.322	.269	.222	.181	.148	.120	.098	.080	• 065	.054	.045	.037	.031	.026	.022
SSURE = I	H2	.000	.001	.004	0.0091	.017	.029	-048	.073	.106	.146	.193	.243	.296	.348	.397	.442	.481	.515	. 543	.566	.584	.599	.611	.621	.629	.635	.640	.644	.647
PRESSI	I	0	N	S	375	0	N	S	-	0	N	5	-	0	N	IO.		0	N	5	-	0	2	5	~	0	N	10		0
•		.200	.200	. 200	0.2019	.203	. 206	.210	.215	.222	.230	.240	.250	.261	.271	.281	.290	.298	.304	.310	.314	.318	.320	.323	.324	.326	.327	.328	.329	.329
20.0 ATM	NH3	9	-	4	0.7885	.778	.762	.739	.707	•666	.616	.558	.496	.432	.369	.310	.257	.211	.172	Ģ	.113	.091	.074	.061	.050	.041	.034	.029	.024	-
11	H2	.000	. 002	.004	9600.0	.017	.031	.050	.077	.111	.153	.201	.252	.306	.358	.407	.452	490	. 523	. 550	.572	.590	.604	.615	.624	.631	.637	.642	.646	649
PRESSURE	Σ	0	N	5	375	0	N	5	-	0	N	5	-	0	2	S		0	N	S	-	0	N	S	-	0	2	5	-	1000

•	~	.200	.200	.200	0.2014	.202	.204	.207	.211	.216	.223	.230	.239	-249	.258	-268	.277	.286	.293	.300	.305	.310	.314	.317	.320	.322	.324	.325	-326	.327	
200.0 ATM.	I	.799	. 798	. 796	0.7918	. 784	. 773	.756	.732	.700	199.	•614	.562	.505	• 446	• 389	.334	.283	.238	.198	.165	.137	.113	*094	*078	• 066	.055	• 046	•039	• 034	
II	HZ	00.	90.	00.	0.0068	O.	.02	.03	• 05	•08	-	•15	•13	•24	•29	.34	•38	•43	• 46	.50	.52	. 55	.57	• 58	99.	.61	*62	•62	• 63	(4)	
PRESSUR	3	0	N	S	375	0	N	S	-	0	N	5	-	0	N	5	-	0	N	5	-	0	N	S	~	0	2	5	-	0	
•	N2	.200	.200	.200	0.2014	-202	.204	.207	-211	.217	.223	.231	.240	.250	.260	.269	.279	.287	.294	.301	.306	* 2 2 3 3 4 5 8	.315	.318	.320	.322	.324	.325	.327	.327	
90.0 ATM.	I	.799	.798	.795	0.7915	.784	.772	.754	.730	1691	.657	609.	555	498	439	381	.326	.275	.231	.192	159	.131	109	.090	.075	.063	.053	.044	.038	.032	
u u	H2	.000	.001	.003	0.0071	.013	.023	.037	.058	.085	.118	.158	. 203	.251	.300	.348	.394	.436	474	.506	.533	.556	.575	.591	.603	.614	.622	.629	.635	•639	
PRESSURE	X	0	N	5	375	0	N	5	-	0	N	5	-	0	N	5	-	0	N	5	-	0	N	5	-	0	N	10	-	0	
•		.20	.20	.20	0.2015	.20	.20	.20	.2I	.21	. 22	.23	. 24	.25	.26	.27	. 28	.28	.29	.30	.30	.33	.31	33	.32	.32	.32	.32	.32	.32	
180.0 ATM.	NH3	.799	.798	795	0.7912	.783	.771	.753	.727	.694	.652	+094	.549	490	.431	.372	.318	.267	.223	.185	.153	.126	.104	.086	.072	.060	.050	.042	.036	.030	
#	7	000	.001	. 003	0.0073	.013	. 024	.039	.060	.088	.122	.163	. 208	.257	.307	.355	.401	.443	.480	.512	. 538	.561	.579	. 594	.606	.616	.624	.631	.636	.640	
PRESSURE	Σ	30	N	S	375	0	N	S	-	0	N	5	-	0	N	5	-	0	2	S	-	0	N	S	-	0	N	5	-	0	

300	C T										
0	7	I		Œ	H2	I	N2	Z	N	m	22
	.000	.799	.20	0	.000	. 799	.200	0	00	.799	.200
N	.001	. 798	. 20	N	.001	. 798	.200	N	00	. 798	.200
S	.003	. 796	.20	5	.003	. 796	.200	S	00	.796	.200
-	.006	.792	.20	-	.006	.792	. 201	-	8	. 792	.201
0	.012	.785	.20	0	.012	.785	.202	0	50	.785	.202
N	.021	.774°	.20	N	.021	.774	. 204	\sim	.02	.775	204
S	.035	.757	\$20	5	.034	.758	.206	5	.03	.759	.206
-	.054	. 734	.21	-	.053	.736	.210		.05	.737	.210
0	.080	.703	.21	0	.078	.706	.215	\circ	.07	. 708	.215
N	.112	.665	• 22	N	.109	*668	.221	N	10	.672	.221
S	.150	e 619	.23	5	.146	.624	.229	S	.14	.628	.228
-	.193	.567	.23	-	.189	.573	.237	-	.18	.578	.237
0	.240	.511	.24	0	.235	.518	.247	0	.23	.523	.246
N	.288	454	25	N	.282	.460	.256	2	.27	.467	.255
S	.336	.396	.26	5	.330	.403	.266	S	.32	.410	.264
شما	.382	341	27	-	.376	.348	.275	1	37	.355	.274
0	.424	.290	28	0	418	.297	.283	0	41	.304	.282
N	.462	. 245	\$29	2	.456	.251	.291	S	.45	,258	.290
S	0.4957	0.2052	0.2991	750	9064-0	0,2113	0.2981	750	0.4857	0.2172	0.2971
-	. 524	1110	330	عصرأ	.519	J 76	.303	-	5	.181	303
0	.548	.142	930	0	,544	.147	.308	0	.54	151*	308
S	. 568	.118	31	N	.564	.122	315	\sim	.56	.126	.312
S	. 584	.098	3	5	.581	.102	.316	S	.57	.105	315
-	.598	.082	E	-	.595	.085	.319	-	.59	.088	.318
0	609.	.068	.32	0	.60T	.071	.321	0	99	.074	.320
3	.618	.057	32	N	.616	090	.323	2	19.	.062	.322
Ŝ	.625	049	* 32	5	.624	.051	.324	S	.62	.053	.324
-	.631	140.	32	100	.630	.043	.326	-	.62	.045	325
	,		1								

HZ NH3 NZ TEMP HZ NH3 NZ TEMP HZ NH3 NZ NZ <th>PRESSURE =</th> <th>240.0 ATM.</th> <th>-</th> <th>PRESSURE</th> <th>URE = 2</th> <th>50.0 ATM</th> <th>•</th> <th>PRESSURE</th> <th>H</th> <th>260.0 AIM</th> <th>•</th>	PRESSURE =	240.0 ATM.	-	PRESSURE	URE = 2	50.0 ATM	•	PRESSURE	H	260.0 AIM	•
0005 0.7994 0.2001 300 0.0005 0.7994 0.2001 300 0.00015 0.7994 0.2002 325 0.0013 0.7985 0.2002 325 0.0013 0.7985 0.2002 325 0.0013 0.7985 0.2002 325 0.0012 0.7985 0.2002 375 0.0012 0.7985 0031 0.7024 0.2002 0.7024 0.7181 0.2002 400 0.0111 0.7867 0.2012 400 0.0118 0.7761 0.2023 400 0.0119 0.7867 0.2023 400 0.0119 0.7867 0.2023 400 0.0119 0.7867 0.2023 400 0.0118 0.7867 0.2033 475 0.00318 0.7871 0.7879 0.0033 475 0.0048 0.7879 0.0034 475 0.0048 0.7879 0.0048 0.7879 0.7879 0.7879 0.0048 0.7879 0.7879 0.7879 0.0048 0.7879 0.0048 0.7879 0.0048 0.7879	2			E	H2	7	N2	T T	HZ	NH3	NZ
0013 0.7844 0.2003 325 0.0013 0.7845 0.2006 355 0.0012 0.7965 0.2006 350 0.0012 0.7965 0.2006 350 0.0012 0.7965 0.0002 0.7965 0.2006 350 0.0013 0.7965 0.0013 0.7013 0.7013 0.7013 400 0.0111 0.7867 0.2012 400 0.0118 0.7767 0.2012 400 0.0119 0.7767 0.2013 425 0.0194 0.7767 0.2013 425 0.0194 0.7767 0.2013 425 0.0194 0.7767 0.2013 425 0.0194 0.7767 0.2013 425 0.0194 0.7767 0.2013 425 0.0194 0.7767 0.2014 475 0.0014 0.7767 0.7167 0.7273 0.0177 0.7173 0.0010 0.0110 0.0010 0.0010 0.0110 0.7767 0.0010 0.0110 0.7767 0.7274 0.7767 0.7274 0.7774 0.7774 0.7774 0.7774	000	0.799	0	0	.000	• 199	.200	\circ	00.	• 799	0.2001
0.00 0.00 <th< td=""><td>9</td><td>0.798</td><td>. 20</td><td>N</td><td>.001</td><td>.798</td><td>.200</td><td>N</td><td>80.</td><td>. 798</td><td>.200</td></th<>	9	0.798	. 20	N	.001	.798	.200	N	80.	. 798	.200
0.0051 0.7927 0.2012 375 0.0059 0.7929 0.2012 375 0.0057 0.7931 0.114 0.7863 0.2023 400 0.0111 0.7867 0.2023 425 0.0199 0.7173 0.126 0.2040 425 0.0114 0.7869 0.2033 425 0.0199 0.7173 0.2033 450 0.0109 0.7173 0.326 0.786 0.2065 475 0.0109 0.7171 0.2083 475 0.0480 0.7173 0.742 0.7110 0.2108 550 0.01016 0.6486 0.2145 500 0.0707 0.7152 1041 0.6751 0.2208 550 0.1711 0.6358 0.2274 500 0.0707 0.7152 11808 0.6321 0.621 0.7410 0.2274 500 0.2145 0.6489 0.2274 0.6589 0.2274 500 0.0214 0.7410 0.2203 500 0.0174 0.7410 0.2203 <	8	0.796	.20	5	.002	• 196	.200	S	9	. 196	.200
0114 0.7863 0.2023 400 0.0111 0.7867 0.2023 400 0.0111 0.7867 0.2023 425 0.0189 0.7773 425 0.0189 0.7761 0.2039 0.7761 0.2039 0.7675 0.0189 0.7773 425 0.0189 0.7767 0.2033 450 0.0318 0.7767 0.2039 475 0.0480 0.7773 475 0.0480 0.7773 475 0.0480 0.7773 475 0.0480 0.7773 475 0.0480 0.7773 0.7771 0.2038 475 0.0480 0.7773 0.0724 0.7731 0.2033 0.2203 0.7628 0.0777 0.7771 0.0727 0.7744 0.775 0.0727 0.0274 0.0727 <t< td=""><td>9.</td><td>0.792</td><td>. 20</td><td>-</td><td>.005</td><td>. 792</td><td>.201</td><td>-</td><td>9</td><td>. 793</td><td>.201</td></t<>	9.	0.792	. 20	-	.005	. 792	.201	-	9	. 793	.201
0199 0.7761 0.2040 425 0.0194 0.7767 0.2039 425 0.0199 0.7771 0.2063 425 0.0189 0.7773 0.0326 0.7669 0.2065 450 0.0318 0.7619 0.2063 450 0.00724 0.7119 0.2068 475 0.0480 0.7724 0.0742 0.7110 0.2088 475 0.0472 0.7111 0.2145 50 0.0474 0.7444 0.0742 0.7131 0.2145 50 0.0104 0.7131 0.2145 50 0.0474 0.7445 0.0994 0.7444 0.0994 </td <td>.01</td> <td>0.786</td> <td>.20</td> <td>0</td> <td>.011</td> <td>.786</td> <td>. 202</td> <td>0</td> <td>£0.</td> <td>. 787</td> <td>.202</td>	.01	0.786	.20	0	.011	.786	. 202	0	£0.	. 787	.202
.0326 0.7609 0.2065 450 0.0318 0.7619 0.2063 450 0.0310 0.7629 .0505 0.7314 0.20492 0.7419 0.2048 475 0.0480 0.7424 .0742 0.7110 0.2148 500 0.07024 0.7141 0.2048 475 0.0494 0.7424 .1042 0.7511 0.6780 0.2203 550 0.0134 0.7624 0.7624 .1399 0.6321 0.2280 550 0.1348 0.6358 0.2203 550 0.1339 0.6391 .1808 0.5830 0.2242 0.7676 0.2744 550 0.1339 0.6391 .2255 0.5264 0.2512 0.4246 0.2244 570 0.1339 0.6391 .2755 0.4168 0.6354 0.2644 0.2244 570 0.1339 0.6394 .3649 0.2639 0.2639 0.3644 0.2719 0.2146 0.3053 .4046 0.2639 <	.01	0.776	.20	N	.019	.776	.203	N	.01	.777	.203
.0505 0.7394 0.2101 475 0.0492 0.7410 0.2098 475 0.0480 0.772 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7724 0.7727 0.7274 0.5274 0.5274 0.5274 0.5274 0.5274 0.5727 0.7274	•03	0.760	.20	S	.031	.761	.206	S	.03	.762	.206
.0742 0.7110 0.2148 500 0.0724 0.7131 0.2145 500 0.0707 0.7152 .1041 0.6751 0.2208 525 0.1016 0.6780 0.2203 525 0.0994 0.6808 .1399 0.67810 0.5784 0.2274 550 0.1339 0.6893 .1808 0.5830 0.2362 575 0.1171 0.5346 0.2274 590 0.1339 0.6893 .2255 0.5294 0.2816 0.2242 600 0.2170 0.5394 0.65918 .2723 0.4733 0.2674 625 0.2674 0.4790 0.2535 625 0.2170 0.5394 0.65918 .3193 0.4168 0.2674 625 0.2674 0.4790 0.2535 625 0.2629 0.4289 .3193 0.4168 0.2674 0.4790 0.2638 625 0.2629 0.4289 .4464 0.2763 0.2684 0.2786 0.4780 0.2789	• 05	0.739	.21	-	•049	,741	.209	-	\$. 742	.209
1041 0.6751 0.2208 525 0.1016 0.6780 0.2203 525 0.0994 0.6808 1399 0.6321 0.2280 550 0.1368 0.6358 0.2274 550 0.1339 0.6393 21808 0.6381 0.2262 0.2362 550 0.1771 0.5875 0.2344 550 0.1735 0.6393 2725 0.6394 0.2451 600 0.2212 0.5346 0.2452 60.1771 0.5346 0.2453 0.6594 0.6394 0.65918 2725 0.4730 0.2460 0.2423 0.2628 650 0.2170 0.5354 0.6759 0.4289 3649 0.3622 0.2730 675 0.3684 0.2719 675 0.4289 0.6759 0.4289 4076 0.3627 0.3684 0.2719 675 0.3624 0.3171 0.3624 0.3749 0.2628 0.5060 0.3684 0.2719 675 0.3741 0.3749 4464	.07	0.711	21	0	.072	.713	.214	0	.07	.715	.214
1399 0.6321 0.2280 550 0.1368 0.6358 0.2274 550 0.1339 0.6393 1808 0.5830 0.2362 575 0.171 0.5875 0.2354 575 0.1735 0.5918 2255 0.5234 0.2421 0.5875 0.2442 600 0.2170 0.5396 2723 0.4733 0.2544 625 0.2674 0.4740 0.2535 625 0.2629 3193 0.4168 0.2639 650 0.3142 0.4236 0.2728 650 0.3093 0.4289 3649 0.3622 0.2815 70 0.4024 0.4179 0.2728 650 0.3034 0.4289 4649 0.2623 0.2893 725 0.4152 0.2719 0.2883 725 0.4289 4464 0.2643 0.2883 725 0.4162 0.2883 725 0.4289 4464 0.2822 0.2883 725 0.4286 0.2883 726 <t< td=""><td>01.</td><td>0.675</td><td>.22</td><td>N</td><td>101.</td><td>.678</td><td>.220</td><td>N</td><td>60.</td><td>.680</td><td>.219</td></t<>	01.	0.675	.22	N	101.	.678	.220	N	60.	.680	.219
1808 0.5830 0.2362 575 0.1771 0.5875 0.2354 575 0.1735 0.5946 2255 0.5294 0.2451 600 0.2212 0.5346 0.2442 600 0.2170 0.5396 2723 0.4733 0.2544 625 0.2674 0.4730 0.2535 625 0.2629 0.4846 3193 0.4168 0.2639 650 0.3142 0.4228 650 0.3093 0.4289 3649 0.3622 0.2730 0.5628 0.365 0.3745 0.3745 4464 0.2643 0.2893 720 0.4762 0.2893 725 0.4415 0.2702 0.2893 725 0.4415 0.2702 0.2893 725 0.4415 0.2702 0.2893 725 0.4415 0.2952 775 0.3546 0.1912 0.2893 725 0.4717 0.2334 4464 0.2623 0.3021 775 0.5065 0.1611 0.3065 0.5283 775<	.13	0.632	.22	5	.136	.635	.227	S	.13	.639	.226
2255 0.5294 0.2451 600 0.2212 0.5346 0.2442 600 0.2170 0.5396 2723 0.4733 0.2544 625 0.2674 0.4790 0.2535 625 0.2629 0.4846 3193 0.4168 0.2639 650 0.3142 0.4230 0.2628 650 0.3093 0.4846 33649 0.3622 0.2730 675 0.3596 0.3684 0.2719 675 0.3745 0.3745 0.3745 4076 0.3109 0.2893 725 0.4024 0.3171 0.2893 725 0.4415 0.2702 0.2883 725 0.4364 0.3744 <	. 18	0.583	23	-	.177	.587	.235	_	11.	.591	.234
2723 0.4733 0.2544 625 0.2674 0.4790 0.2535 625 0.2629 0.4846 3193 0.4168 0.2639 650 0.3142 0.4230 0.2628 650 0.3093 0.4289 3649 0.3622 0.2730 675 0.3644 0.2719 675 0.3546 0.3745 4076 0.3103 0.4024 0.3171 0.2805 700 0.3974 0.3745 4076 0.2893 725 0.4415 0.2702 0.2883 725 0.4367 0.2760 4464 0.2843 725 0.4415 0.2702 0.2883 725 0.4367 0.2760 4809 0.2229 0.2952 750 0.4762 0.2285 0.2952 750 0.4371 0.2339 5510 0.1871 0.3023 0.4189 0.3183 825 0.5284 0.1611 0.3065 0.1094 0.3165 0.1180 0.5145 0.1659 0.1659 0.1659 0.16	. 22	0.529	.24	0	.221	.534	.244	0	.21	.539	.243
3193 0.4168 0.2639 650 0.3142 0.4230 0.2628 650 0.3093 0.4289 3649 0.3622 0.2730 675 0.3646 0.2719 675 0.3546 0.3745 4076 0.2181 700 0.4024 0.3171 0.2805 700 0.3974 0.3745 4464 0.2263 0.2893 725 0.4415 0.2702 0.2883 725 0.4367 0.3746 4809 0.2229 0.2962 750 0.4717 0.2339 4809 0.2229 0.2962 750 0.4717 0.2339 5362 0.1871 0.3021 775 0.4717 0.2339 5362 0.1871 0.3021 775 0.4717 0.2339 5576 0.1879 0.3065 0.3065 0.3165 0.3065 0.4717 0.3287 0.1656 5575 0.1094 0.3145 0.3145 825 0.5587 0.1656 5575 <t< td=""><td>.27</td><td>0-473</td><td>25</td><td>N</td><td>-267</td><td>.479</td><td>.253</td><td>O</td><td>.26</td><td>.484</td><td>.252</td></t<>	.27	0-473	25	N	-267	.479	.253	O	.26	.484	.252
.3649 0.35622 0.2730 675 0.3596 0.3684 0.2719 675 0.3546 0.3745 .4076 0.2815 700 0.4024 0.3171 0.2805 700 0.3974 0.3231 .4644 0.2643 0.2893 725 0.4415 0.2702 0.2883 725 0.4367 0.2760 .4809 0.2229 0.2962 750 0.4717 0.2760 0.2760 0.7760 0.2760 0.	.31	0.416	.26	5	.314	.423	.262	5	30	.428	.261
40760.31090.28157000.40240.31710.28057000.39740.323144640.26430.28937250.44150.27020.28837250.43670.276048090.22290.29627500.47170.27397500.47170.233948090.22290.29627750.47620.22850.29527750.47170.233953620.18710.30217750.50650.19220.30137750.650240.197153620.18750.013490.31648250.55420.16110.30658000.52870.165655760.10940.31518500.55420.11340.31458500.55900.116559020.09170.31808750.58760.09490.31458750.58760.094960240.07710.32059000.60140.06750.32009000.59790.069861240.06510.32259250.61890.05239250.60850.069862740.04710.32560.06730.32529750.06360.063663300.04040.32640.04190.326310000.63050.0434	36	0.362	.27	-	.359	.368	.271	-	.35	.374	.270
4464 0.2643 0.2893 725 0.4415 0.2702 0.2883 725 0.4415 0.2702 0.2952 750 0.4717 0.2339 4809 0.2229 0.2962 750 0.4762 0.2285 0.2952 750 0.4717 0.2339 5108 0.1871 0.3021 775 0.5065 0.1922 0.3013 775 0.5024 0.1971 0.2339 5362 0.1871 0.3065 0.3065 0.1611 0.3065 800 0.5287 0.1971 5576 0.1308 0.3151 850 0.5542 0.1130 0.3145 850 0.5509 0.1138 5575 0.1094 0.3151 875 0.5542 0.1130 0.3145 850 0.5509 0.1165 5502 0.0917 0.3180 875 0.5876 0.0949 0.3145 875 0.5509 0.1165 6024 0.0771 0.3225 925 0.6104 0.0675 0.3238 950 <td>64.</td> <td>0.310</td> <td>28</td> <td>0</td> <td>.402</td> <td>.317</td> <td>.280</td> <td>\circ</td> <td>•39</td> <td>.323</td> <td>.279</td>	64.	0.310	28	0	.402	.317	.280	\circ	•39	.323	.279
.4809 0.2229 0.2962 750 0.4762 0.2285 0.2952 750 0.4717 0.2339 .5108 0.1871 0.3021 775 0.5065 0.1922 0.3013 775 0.5024 0.1971 .5362 0.1565 0.3072 800 0.5324 0.1611 0.3065 800 0.5287 0.1971 .5362 0.3115 825 0.5542 0.1611 0.3168 825 0.5587 0.1656 .5755 0.1094 0.3151 850 0.5572 0.1130 0.3145 850 0.5599 0.1165 .5902 0.0917 0.3180 875 0.5876 0.0949 0.3175 875 0.5876 0.1165 .6024 0.0971 0.3205 900 0.6001 0.0799 0.3200 900 0.5979 0.0678 .6124 0.0651 0.3225 925 0.6189 0.0573 0.3238 950 0.6189 0.0678 0.9328 975 0.625	4.	0.264	.28	N	.441	.270	.288	M	.43	.276	.287
-5108 0.1871 0.3021 775 0.5065 0.1922 0.3065 800 0.5324 0.1611 0.3065 800 0.5287 0.1656 .5362 0.1565 0.3065 800 0.5324 0.1611 0.3065 800 0.5287 0.1656 .575 0.1094 0.3115 825 0.5542 0.1130 0.3108 825 0.5509 0.1389 .5755 0.1094 0.3151 850 0.5725 0.1130 0.3145 850 0.5596 0.1165 .5902 0.0917 0.3180 875 0.5876 0.0949 0.3175 875 0.5891 0.0979 .6024 0.0771 0.3225 925 0.6104 0.0675 0.3221 925 0.6104 0.0675 0.3231 925 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0419 0.3253 1000 0.6305 0.0434 0.3253 1000 0.6305 0.0434	.48	0.222	.29	5	.476	.228	.295	S	14.	.233	.294
.5362 0.1565 0.3072 800 0.5324 0.1611 0.3065 800 0.5287 0.1656 .5576 0.1308 0.3165 825 0.5542 0.1349 0.3108 825 0.5509 0.1389 .5755 0.1094 0.3151 850 0.5725 0.1130 0.3145 850 0.5696 0.1165 .5902 0.0917 0.3180 875 0.5876 0.0949 0.3175 875 0.1659 .6024 0.0971 0.3205 900 0.6001 0.0799 0.3200 900 0.5851 0.0979 .6124 0.0651 0.3225 925 0.6104 0.0675 0.3221 925 0.6104 0.0675 0.3221 925 0.6109 .6204 0.0553 0.3241 950 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0419 0.3255 975 0.6259 0.0419 0.3253 1000 0.6305 0.04	.51	0.187	30	-	.506	.192	.301	-	.50	.197	.300
.5576 0.1308 0.3115 825 0.5542 0.1349 0.3108 825 0.5509 0.1389 .5755 0.1094 0.3151 850 0.5725 0.1130 0.3145 850 0.5696 0.1165 .5902 0.0917 0.3180 875 0.5876 0.0949 0.3175 875 0.5651 0.0979 .6024 0.0771 0.3205 900 0.6001 0.0799 0.3200 900 0.5979 0.0825 .6124 0.0651 0.3225 925 0.6104 0.0675 0.3221 925 0.6085 0.0698 .6206 0.0553 0.3241 950 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0471 0.3255 975 0.6259 0.0489 0.3255 975 0.6259 0.0419 0.3253 1000 0.6305 0.0434	.53	0.156	.30	0	.532	.161	.306	\Box	.52	.165	.305
.5755 0.1094 0.3151 850 0.5725 0.1130 0.3145 850 0.5866 0.1165 .5902 0.0917 0.3180 875 0.5876 0.0949 0.3175 875 0.5851 0.0979 .6024 0.0771 0.3205 900 0.6001 0.0799 0.3200 900 0.5979 0.0825 .6124 0.0651 0.3225 925 0.6104 0.0675 0.3221 925 0.6098 .6206 0.0553 0.3241 950 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0471 0.3255 975 0.6259 0.0489 0.3252 975 0.6245 0.0506 .6330 0.0404 0.3256 1000 0.6317 0.0419 0.32563 1000 0.6305 0.0434	. 55	0.130	31	2	.554	.134	.310	O.	.55	• 138	.310
.5902 0.0917 0.3180 875 0.5876 0.0949 0.3175 875 0.5851 0.0979 .6024 0.0771 0.3205 900 0.6001 0.0799 0.3200 900 0.5979 0.0825 .6124 0.0651 0.3225 925 0.6104 0.0675 0.3221 925 0.6085 0.0698 .6206 0.0553 0.3241 950 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0471 0.3255 975 0.6259 0.0489 0.3252 975 0.6245 0.0506 .6330 0.0404 0.3266 1000 0.6317 0.0419 0.3263 1000 0.6305 0.0434	.57	0.109	33	5	.572	.113	.314	S.	.56	1116	.313
.6024 0.0771 0.3205 900 0.6001 0.0799 0.3200 900 0.5979 0.0825 .6124 0.0651 0.3225 925 0.6104 0.0675 0.3221 925 0.6085 0.0698 .6206 0.0553 0.3241 950 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0471 0.3255 975 0.6259 0.0489 0.3252 975 0.6245 0.0506 .6330 0.0404 0.3266 1000 0.6317 0.0419 0.3263 1000 0.6305 0.0434	539	0.091	31	-	.587	*60	.317	-	.58	.097	.317
.6124 0.0651 0.3225 925 0.6104 0.0675 0.3221 925 0.6085 0.0698 .6206 0.0553 0.3241 950 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0471 0.3255 975 0.6259 0.0489 0.3252 975 0.6245 0.0506 .6330 0.0404 0.3266 1000 0.6317 0.0419 0.3263 1000 0.6305 0.0434	• 60	0.077	32	0	009*	.079	.320	0	.59	.082	• 318
.6206 0.0553 0.3241 950 0.6189 0.0573 0.3238 950 0.6172 0.0593 .6274 0.0471 0.3255 975 0.6259 0.0489 0.3252 975 0.6245 0.0506 .6330 0.0404 0.3266 1000 0.6317 0.0419 0.3263 1000 0.6305 0.0434	.61	0.065	.32	S	.610	.067	.322	\sim	.60	• 069	.321
.6274 0.0471 0.3255 975 0.6259 0.0489 0.3252 975 0.6245 0.0506 .6330 0.0404 0.3266 1000 0.6317 0.0419 0.3263 1000 0.6305 0.0434	.62	0.055	32	5	.618	.057	.323	5	.61	• 059	.323
.6330 0.0404 0.3266 1000 0.6317 0.0419 0.3263 1000 0.6305 0.0434	. 62	0.047	32	-	.625	.048	.325	-	.62	.050	.324
	. 63	0.040	.32	8	.631	.041	.326	$\boldsymbol{\circ}$.63	•043	.326

8	270.0 ATM.	•	PRESSURE	n	280.0 ATM.		w	II	290.0 ATM.	
, ,,,,,,	I		Z	H2	NH3	N2	Σ	H2	NH3	S
ö	66	~	0	000	-	.200	0	.000	. 199	• 20
0	198	3	S	001		.200	N	.001	. 798	•20
o	196	3	5	002		.200	S	.002	.796	.20
Ö	193	~	-	005	-	.201	-	.005	.793	• 20
0	787	3	0	010	0	.202	0	.010	. 787	.20
O	111	?	N	018	-	.203	S	.017	. 778	• 20
Ö	763	3	5	029		.205	5	.028	.765	.20
0	743	4	-	045		.209	-	.044	.746	.20
0	717	2	0	190	-	.213	0	• 066	.720	.21
0	683	S	N	095	*	.219	N	.093	.688	.21
0	642	N	S	128		.225	5	.126	.648	.22
0	595	?	-	166	, so	.233	-	.163	.603	.23
0	.5443	0.2426	909	0.2093	0.5488	0.2419	009	0.2058	0.5531	0.2412
0	489	2	N	254	4	.250	N	.250	665.	.25
0	434	S	5	300	4	.260	5	.295	.445	.25
~	380	3	معيرا	345		,269	-	.340	.391	• 26
0	328	2	0	387	4	.277	0	.383	.339	.27
•	281	2	2	427	S	.285	N	.423	.292	.28
Q.	239	7	S	463	C)	.292	S	.458	.249	.29
_	202	2	-	464	CA	.298	-	490	.211	.29
J	170	·W	0	521	7	.304	0	.517	.178	.30
	142	S.	N	544	the state of	.308	~	.541	.150	8
•	120	W.	S	563	,; #	.312	S	.561	.126	.31
<u>_</u>	101	S.	-	580	-	.316	-	.577	.106	.31
0	085	4	0	593	9	.318	0	.591	060°	.31
_	072	ୁ	N	409	9	.320	N	.602	.076	.32
-	061	3	5	613	9	.322	S	.612	.065	.32
_	N	e U	-	621	Ç	.324	-	.620	.055	.32
O	.044	(A)	9	628	9	.325	0	.626	.047	.32

•	N2	- 200	.200	.200	0.2010	.201	.203	.205	.208	.212	.217	.223	.231	.239	.247	.256	.265	.274	.282	.289	.295	.301	.306	.310	.314	.317	.319	.321	.323	.324
320.0 ATM.	m	• 799	• 798	. 797	0.7940	• 788	. 780	.767	.749	.725	+69*	• 656	.613	.564	.512	.459	• 406	.355	.307	.263	.224	190	191.	.136	.115	160.	.083	.071	090-	.052
H	H2	000	.001	.002	0.0050	• 000	•016	.027	•042	.062	.088	.119	.155	.196	.239	.283	.328	.370	.410	.447	6140	.507	.532	.552	.570	.585	.597	.607	619	•623
PRESSURE	2	0	N	5	375	0	N	S	-	0	N	5	-	0	N	5	-	0	N	in	-	0	N	5	-	0	N	5	-	
	N2	.200	.200	.200	0.2010	.201	.203	.205	.208	-212	.217	-224	.231	.239	.248	.257	.266	.275	.282	.290	.296	.302	.307	.311	.314	.317	.319	.321	.323	.324
310.0 ATM.	エ	£661°	. 798	197	0.7939	. 788	.779	. 766	.748	.723	.692	.654	.610	.561	. 508	.454	.401	.350	.302	.258	.220	. 186	.157	.133	.112	.095	.081	•069	.059	.050
II W	H2	000	.001	.002	0.0051	•000	.016	.027	.043	.063	• 089	.121	.158	.199	.242	.287	.332	.374	.414	.451	.483	.511	.535	.555	.572	.587	.599	609	.617	. 624
PRESSUR	Σ	0	N	5	375	0	N	5	-	0	N	5	-	0	N	S	-	0	N	5	-	0	N	S	-	0	N	5	-	0
.•	N2	N	3	~	0.2010	8	2	Š	3	7	2	7	N	~	3	?	7	~	~	2	~	.0	S	3	6	3	(U)	4	u,	(C)
300.0 ATM.	NH3				0.7937	-			-	-	9	9	•	r.	*	4	6	(F)	3	2	4	7	7			0	0	0	0.0575	0
H	2	.000	.001	.002	0.0052	.009	. 017	.028	.043	.064	.091	.123	.161	202	. 246	. 291	.336	.379	.418	454	.486	.514	.538	.558	.575	.589	.600	.610	.618	•625
PRESSURE	TEMP	300	325	350	375	400	425	450	475	500	525	550	575	9	625	650	675	700	725	750	775	800	825	850	875	900	925	950	975	1000

	N2	.200	-200	.200	.200	.201	. 203	-205	.208	.211	.216	.222	.229	.237	.245	.254	.263	.271	.279	.287	.293	.299	.304	.309	.312	.315	.318	.320	0.3224	•323
350.0 ATM		. 799	. 798	. 797	.794	.789	. 781	.769	.752	.729	669.	.664	.622	.575	.524	.472	.419	.368	.320	.276	.236	.202	.171	.145	.123	.105	.089	.076	0.0658	• 056
SSURE = 3	H2	00.	9	8	8	00.	50	.02	.03	• 05	• 08	11.	.14	8	.22	.27	.31	.35	.39	•43	•46	640	.52	• 54	.56	.57	• 59	• 60	0.6118	19.
PRESSU	2.	0	S	S		O	S	S	-	0	N	S	-	O	~	S	-	0	V	S	and a	\circ	N	£.	-	\circ	N	S.	975	0
		. 200	.200	.200	.201	.201	.203	.205	.208	.212	.217	.223	.230	.238	.246	.255	.264	.272	.280	.287	.294	.300	.305	.309	.313	.316	.318	.320	0.3226	.324
340.0 ATM.	I	. 799	. 798	.797	.794	.789	.781	.768	.751	. 728	• 698	199.	.619	.571	.520	.468	.415	.364	.316	.272	.232	.198	.168	.142	.121	. 102	.087	.074	0.0642	•055
SSURE = 34	H2	000	.001	.002	.004	600.	.015	.026	.040	.060	.084	-115	.150	.190	.232	.276	.320	.363	.403	.439	.472	.501	.526	.547	.565	.581	.593	. 604	0.6132	.620
PRESSI	E	0	2	5		0	2	5		0	S	5	-	0	N	5	-	0	N	5	-	0	2	5	-	0	N	5	975	0
		.200	. 200	.200	. 201	.201	. 203	.205	.208	.212	.217	.223	.230	.238	.247	.256	.264	.273	.281	.288	-295	.300	.305	.310	.313	.316	.319	-321	0.3229	.324
330.0 ATM.	I	.79	. 79	• 79	.79	.78	. 78	.76	. 75	.72	69.	.65	.61	• 56	.51	.46	.41	. 35	.31	.26	*22	• 19	. 16	•13	11.	• 10	.08	.07	0.0625	• 05
П	귚	.000	.001	.002	.004	.009	.016	.026	.041	.061	.086	.117	.153	• 19	.235	.280	. 324	.366	. 406	• 443	.476	. 504	. 529	.550	. 568	.583	.595	.605	0.6145	•621
PRESSURE	T.	0	2	5	-	0	2	5	-	0	2	5	-	0	S	5	-	0	N	S	-	0	2	5	_	0	2	5	975	0

PRESS	SSURE = 3	360.0 ATM	<u>.•</u>	PRESS	SSURE = 3	370.0 ATM.	<u>.</u>	PRESSI	SSURE = 3	380.0 ATM.	"\
2	42	SES		TEMP	H2	I		X	H2	NETS	N2
30	8	.79	.20	300	.000	.79	.200	0	00.	• 79	.200
N	00.	.79	. 20	325	.001	.79	.200	S	8	. 79	.200
in	80.	• 79	.20	350	.002	51.	.200	5	9	• 79	.200
375	0.0046	0.7944	0.2009	375	0.0045	0.7945	0.2009	375	0.0045	0.7946	0.2009
0	00.	.78	. 20	400	.008	• 78	.201	0	00.	. 78	. 201
N	10.	. 78	. 20	425	.015	.78	. 203	2	.01	. 78	-202
5	.02	.76	.20	450	.024	1	.204	5	.02	17.	.204
-	8	.75	. 20	475	.038	. 75	.207	-	.03	• 75	.207
0	.05	.73	.21	200	.057	-73	.211	0	• 05	. 73	.211
N	80	.70	.21	525	.080	.70	.216	2	.07	.70	.215
5	1	• 66	.22	550	.109	• 66	.221	5	.10	.67	.221
-	14	.62	.22	575	.143	.62	.228	-	•14	• 62	.228
0	.18	.57	.23	009	.182	.58	.236	0	.17	. 58	.235
N	.22	.52	*24	625	.223	ES.	.244	2	.22	.53	-244
S	.26	14.	.25	650	.266	*48	.253	5	.26	• 48	-252
~	31	. 42	. 26	675	.309	•42	-262	-	.30	.43	.261
0	35	.37	.27	700	.352	3	.270	0	.34	•38	-269
N	9	.32	*27	725	.392	.32	.278	S	.38	• 333	.277
in	43	.28	.28	750	.429	.28	.285	5	.42	• 28	.285
-	46	.24	.29	775	.462	.24	.292	-	.45	• 24	.291
0	49	. 20	.29	800	.492	.20	.298	0	. 48	.21	.297
N	.52	.17	.30	825	.518	11.	.303	2	.51	18	303
10	54	14	.30	850	.540	.15	.308	5	.53	• 15	.307
~	.56	.12	.31	875	.559	12	.311		.55	E1 •	*311
Ö	.57	.10	.31	900	.575	.11	.315	0	.57	.11	.314
N	.59	• 09	.31	925	.588	60.	.317	N	.58	• 09	17
I	.60	.07	.32	950	.599	.08	.319	5	.59	80.	.319
-	.61	•06	.32	975	609	0.0691	.321	~	• 60	6	
0	.61	.05	.32	1000	.617	53	•323	1000	0.6158	90.	0.3232

3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3 3	1		!	1	1	;		1	()
NH3 N2	N2	TEMP	1 2	EE 3	N2	TEMP	1 5	E P	2
.7996 0.200	.200	0	.000	• 799	. 200	0	.000	. 799	.200
.7989 0.200	.200	2	.000	. 798	.200	2	000	.798	.200
.7974 0.200	.200	5	.002	167.	.200	5	.002	. 797	.200
.7947 0.200	.200	-	.004	.794	.200	-	•004	. 794	.200
0.7901 0.2017	. 201	400	0.0081	0.7902	0.2016	0	0.0080	0.7904	0.2016
.7826 0.202	. 202	2	.014	.782	.202	2	.014	. 783	.202
.7713 0.204	.204	5	.023	.771	.204	5	.023	.772	.204
.7554 0.207	.207		.036	.756	.207	-	• 036	.756	.207
•7338 0.211	.211	0	.054	.734	.210	0	.053	. 735	.210
.7061 0.215	.215	2	.077	. 707	.215	2	.075	.708	.215
.6722 0.221	.221	5	.105	.674	.221	5	.103	•675	.220
.6323 0.227	.227		.137	.634	.227	-	.136	.636	.227
.5873 0.235	.235	0	.174	.590	.235	0	.172	.592	.234
.5386 0.243	.243	N	.215	. 541	.243	2	.212	.544	.245
.4875 0.252	.252	5	.257	.491	.251	5	.254	*65*	.250
.4360 0.260	.260	-	.300	°439	.260	~	.297	.443	.259
.3854 0.269	.269	0	.342	.389	.268	0	.339	.393	.267
.3373 0.277	.277	2	.382	.341	.276	2	.379	.345	.275
.2927 0.284	.284	5	.419	*296	.283	S	.416	.300	.283
.2522 0.291	. 291	-	.453	*255	. 290	-	.450	•259	.290
.2162 0.297	.297	0	.483	.219	.296	0	.480	-222	.296
.1847 0.302	.302	2	.510	.187	.302	S	.507	190	.301
1575 0.307	.307	5	. 533	.160	.306	5	.530	.163	.306
.1343 0.310	.310	-	.552	.136	.310	-	.550	.139	.310
.1146 0.314	.314	0	.569	.116	.313	0	.567	.119	.313
.0980 0.317	.317	N	.583	.100	.316	2	.581	.102	.316
.0840 0.319	.319	5	.595	.085	.319	5	.593	.087	.318
.0723 0.321	.321	-	.605	.073	.32I	-	03	.075	2
.0624 0.322	.322	0	.613	• 063	.322	1000	612	.065	2

<u>.</u>	N2	• 20	.20	.20	0.2008	• 20	• 20	• 20	.20	.21	.21	.27	•22	.23	• 24	•24	• 25	• 26	.27	.28	•28	• 29	•30	.30	•30	.31	.31	.31	.32	-32
440.0 ATM	I	• 799	• 799	.797	0.7951	. 790	.783	.773	. 758	.738	.712	. 680	•643	009.	.553	.504	• 453	.403	.355	.310	*269	.232	199	.171	.146	.125	.107	.092	.080	• 069
ESSURE = 4	H2	00.	9	80.	0.0041	900	.01	.02	• 03	.05	.07	60.	•13	• 16	•20	.24	• 28	.33	.37	.40	• 44	.47	.50	.52	.54	• 56	.57	. 58	.59	99.
PRESS	2	\circ	S	S	375	\circ	3	5	-	0	2	S		0	2	S	-	0	N	5	-	0	2	5	-	0	2	S	-	0
•	N N	.200	.200	.200	0.2008	.201	.202	.204	.207	.210	.214	.220	.226	.233	.241	.249	.258	.266	.274	.282	.288	.295	.300	.305	.309	.312	.315	.318	.320	.322
430.0 ATM.	NH3	. 79	.79	• 79	0.7951	.79	.78	.77	• 75	•73	.71	.67	• 64	.59	.55	. 50	.45	.40	35	.30	• 26	•22	5T *	• 16	.14	.12	.10	• 09	.07	• 06
SSURE = 4	H2	.000	000	.002	0.0041	.007	.013	.022	.034	.051	.073	.100	.132	.168	.207	.249	.291	.333	.373	.410	.444	.475	.502	.526	.546	.563	.578	.590	.601	•610
PRESS	Œ	0	N	5	375	0	N	S	-	0	N	5	~	0	N	5	-	\circ	N	in	-	\circ	N	5		0	N	S	_	0
•		. 200	.200	.200	0.2008	.201	.202	.204	.207	.210	.215	.220	.226	.234	. 242	.250	.258	.267	.275	.282	.289	.295	.301	.305	.309	.313	.316	.318	.320	.322
20.0 ATM.		. 79	• 79	.79	0.7950	. 79	.78	.77	. 75	.73	.71	19.	.63	.59	.54	64.	44.	96.	3.4	.30	.26	.22	• 19	• 16	.14	.12	100	.08	.07	S
11	H2	000	.000	. 002	0.0042	.007	.013	.022	.035	.052	.074	. 102	.134	.170	.210	.251	.294	.336	.376	.413	.447	.478	. 505	. 528	. 548	.565	.580	. 592	.602	.611
PRESSURE	Z	0	2	S	375	0	S	5	-	0	S	S		0	N	S	-	0	N	S	-	0	N	S	-	0	N	5	-	0

•	N2	.20	.20	• 20	.20	.20	• 20	.20	• 20	.20	.21	.21	.22	.23	0.2398	•24	.25	•26	.27	•27	.28	•29	•29	•30	.30	.31	3	.31	33	2
70.0 ATM.	I	• 799	. 199	197	. 795	. 791	.784	* 774	.760	.741	. 716	.685	.648	.607	56	.513	.463	.413	•366	.321	.279	.241	.208	.179	.153	.132	.113	.097	.084	.073
7	H2	0000	000	.001	.003	.007	.012	.021	.033	.049	.069	.095	.126	.160	0.1988	.239	.280	.321	.361	.399	.433	.465	.492	.517	.538	,556	571	.585	.596	* 605
PRESSURE	Σ	300	2	3	-	0	N	5	-	0	S	S	-	0	2	S	-	0	2	Ś	-	0	N	5	-	\circ	2	S	-	0
		.200	.200	.200	.200	.201	.202	.204	.206	.209	.214	.219	.225	.232		.248	.256	.264	.272	.280	.287	.293	.299	.303	.308	.311	314	317	.319	.321
460.0 ATM.	NH3	. 799	66L°	797	. 795	.791	. 784	,774	.759	.740	215	.683	.646	,604	0.5588	.510	.460	.410	.362	.317	.276	*238	.205	.176	151	.130		960	.083	.072
SSURE = 4	H2	0000	0000	.001	.003	.007	.013	.021	.033	.049	.070	960.	1210	.162	S	241	.283	.324	,364	.40I	.436	1667	.495	.519	.540	.558	573	.586	. 597	• 606
PRESSI	Œ	300	2	5	-	0	N	5	-	0	S	5	-	0	S	5	-	0	2	S	Sept.	\circ	N	S	-	\bigcirc	N	S	and a	0
à	N N	.20	.20	200	.20	.20	20	20	.20	7	22	7	.22	23	4	24	25	.26	120	2.8	.28	.29	.29	30	.30	31	3	.33	33	(L)
450.0 ATM.		.799	.799	. 797	.795	.791	. 784	.773	.759	.739	E113	.682	. 645	.602		.507	.457	.407	.359	.314	.273	,235	.202	174	64T°	.128	109	460°	.081	.070
#	72	.000	.000	.001	.004	.007	.013	.021	. 033	.050	.071	.098	.129	164	S	.244	.285	.327	.367	.404	.439	.470	. 497	.521	. 542	.560	575	.587	.598	607
PRESSURE	Z	300	2	S	-	0	2	S	-	0	N	S	-	0	N	5	-	0	N	S	-	0	S	S	-	0	S	5		O

PRESS	ESSURE = 4	480.0 ATM.	•	PRESSURE	#	490.0 ATM.	•	PRESSURE	H :	500.0 ATM.	
Σ	7	NH3	N2	Σ	H2	I		Σ	N	NH3	
0	000	* 79	.200	0	.000	661 *	,200	0	000	.799	• 200
S	.000	.79	.200	N	.000	. 199	.200	2	.000	* 799	. 200
5	.001	.79	.200	5	.001	197	.200	5	.001	197	. 200
-	.003	. 79	-200		.003	. 795	.200	-	• 003	. 795	.200
0	.007	979	.201	0	.007	.791	.201	\bigcirc	.007	161.	.201
N	-012	.78	.202	S	.012	. 785	.202	N	0175	. 785	.202
5	.020	.77	.204	S	.020	• 775	. 204	5	.020	• 775	.204
-	.032	. 76	.206	-	.032	.761	.206	-	.031	* 761	• 206
0	.048	.74	.209	0	.047	.742	.209	0	• 047	.743	.209
N	.069	.71	.213	N	.068	.718	.213	2	• 067	.719	.213
5	. 094	* 68	.218	5	.093	.688	.218	S	.092	• 689	.218
-	.124	.65	. 224	-	.123	.652	.224	-	.121	.653	.224
0	.158	.60	.231	0	.157	.611	.23I	0	,155	.613	.231
N	196	.56	.239	N	.194	.566	.238	2	.192	.568	.238
S	.236	.51	.247	5	.234	.518	.246	5	.232	.521	-246
-	.278	46	.255	-	.275	.469	.255	-	.273	.472	.254
0	.319	. 4]	. 263	0	.316	.420	. 263	0	.314	.423	.262
2	.358	36	.271	N	,356	.372	.271	3	.353	.375	.270
S	.396	.32	.279	5	.393	.327	.278	5	.391	• 330	.278
-	.43I	• 28	.286	-	.428	.285	.285	معرا	•426	.288	.285
0	.462	.24	-292	0	.460	.247	.292	0	450 J	.250	.291
N	.490	.21	.298	2	.488	.214	.297	S	•486	,216	. 297
5	.515	.18	.303	S	.513	.184	.302	S	.511	.186	.302
-	.536	.15	.307		.534	.158	•306	شما	. 532	•160	.306
0	.554	.13	.310	0	.553	.136	.310	0	. 551	.138	.310
2	.570	.11	.314	2	.568	.1117	.313	2	.567	•119	.313
S	. 583	60.	.316	5	.582	.101	.316	S	.580	.102	.316
975	S	0.0861	0.3190	975	0.5937	0.0876	0.3187	975	0.5924	0.0891	0.3185
0	• 604	.07	.320	0	° 603	• 076	-320	0	• 602	.077	.320

	N2	200	200	200	200	201	207	203	206	507	773	217	223	230	237	245	253	261	269	.276	283	0.2902	29	30	30	30	315	315	31	بن اسم	
530.0 ATM.	NH3	199	199	161	795	791	785	776	, 763	745	, 122	693	658	619	575	528	480	43	384	333	.297	ഹ	.224	194	• 16	.144	.124	2	60	.08]	
.11	H2	000	000	001	003	900	011	019	030	24.5	065	980	111	150	181	.226	.266	306	346	383	.418	0.4508	41	504	.52	.546	56.	.570	.588	.59	
PRESSURE	TEMP	300	325	350	375	400	425	450	415	200	525	550	575	909	625	650	675	700	725	750	775	800	825	850	875	006	925	950	975	1000	
_	~.	200	200	200	200	201	202	204	206	503	213	,218	,223	,230	.237	245	253	.261	269	.277	284	0.2906	,296	301	305	306	315	316	318	,32(
20.0 ATM.	-	799	799	797	795	791	785	776	.762	744	721	691	657	617	573	526	47	429	381	336	294	0.2562	.223	191	.165	142	122	100	092	.070	•
SSURE = 57	H2	00	8	9	8	8	9	2	E	9	9	9	=	-	18	2	2	3	J.	3	4	0.4531	*	N	S	5	S	5	Š	9	
PRESSI	-	30	~ .		-	_	~1	•		~	C)	10	700			L LET	-			11.7	1	800		41	-	_		**	. , ,	_)
•	N2	200	200	200	200	201	202	204	206	209	213	218	224	230	236	246	254	262	27(27	2	0.2911	. 29	30	30	300	3	7	~	200	Į.
510.0 ATM.		7	6	79	79	7	7	7	76	7	-	9	3	Ģ	1	ľ	4	4	6	. (Š	0.2534	~	-	-	-	,		Š	9	2
W	4	9	8		č	5 8	Ö	6	ć	9	90	ě	-	4	-	1 5	,	Ç	י ר	'n	•	0.4555	3	Š	, ic		, E	, ic	` ŭ	•	ō
PRESSURE	-	-			٠.	_	• ~				. ^	, ,,			3 (u 16		• •	36	u u		200		, ,	3 7 7			uu		- 6	~

PRESS	SSURE = 5	540.0 ATM.	<u></u>	PRESS	SSURE = 5	50.0 ATM	•	PRESSI	SSURE = 5	560.0 ATM.	•
X	7	NH3		E	H2	I	N2	2	H2	NH3	N2
300		0.7996	C	300	000	.799	.200	300	000	0.7997	• 20
N		0.7991	S	S	.000	.799	.200	N	000.	29	.20
S	•	0.7979	CA.	S	.001	*798	200	S	.001	6	• 20
-	0.0035	0.7957		-	0.0035	9	0.2007	-	0.0035	0.7958	0.2007
0	•	0.7920	(4	0	.006	. 792	.201	O	•000	79	•20
N	•	0.7859	C)	N	.011	.786	.202	O	.011	78	• 20
S		0.7768	CVI	5	.019	177	.203	S	.018	7	.20
		0.7637	C	-	.029	.764	.206	-	.029	20	• 20
0		0.7460	CA.	0	.044	.746	.208	0	.044	7	• 20
N		0.7229	CAT.	N	.063	.723	.212	2	.062	72	.21
5		0.6943	4	5	.087	.695	.217	S	.086	63	.23
-		0.6601	N	-	.115	.661	.223	-	,114	99	•22
0		0.6209	4	0	.147	.622	.229	0	.146	2	*22
N		0.5775	€	N	.183	.579	.236	N	.182	58	• 23
5		0.5311	CA	5	.222	.533	.244	5	.220	33	•24
~		0.4830	17	-	.262	. 485	.252	-	.259	8	.25
0		0.4347	7	0	.302	.437	.260	0	.299	4	• 26
N			7	3	.341	.390	.268	N	.339	33	•26
S		0.3423	S	5	.379	.345	.275	5	.376	4	.21
-			(7)	-	.414	.302	282	~	.411	8	• 28
0			CA	0	.446	,264	.289	0	444	20	• 28
N			S	N	.475	.229	.295	S	.473	2	•28
S				5	.501	198	300	5	•499	2	.29
-			e.	-	.523	.171	.304	-	.521	_	.30
0			C.	0	.545	.148	.308	0	.541	5	• 30
N		0.1265	123	N	.559	.128	311	2	.558	2	£.
5		0.1094	EJ.	5	.574	111	.314	S	.572	Ξ	.31
-			(1)		.586	• 096	317	-	. 585	8	33
0		0.0825	th.	0	0.5969	.083	0.3194	1000	• 595	0.0850	0.3192

PRESSURE = 590.0 ATM.	MP H2 NH3	00 0.0003 0.7997 0.20	25 0.0007 0.7991 0.20	50 0.0016 0.7980 0.20	75 0.0033 0.7960 0.20	00 0.0063 0.7924 0.20	25 0.0111 0.7867 0.20	50 0.0182 0.7781 0.20	75 0.0285 0.7657 0.20	00 0.0426 0.7489 0.20	25 0.0608 0.7270 0.21	50 0.0836 0.6997 0.21	75 0.1108 0.6670 0.22	00 0.1422 0.6293 0.22	25 0.1772 0.5874 0.23	50 0.2147 0.5423 0.24	75 0.2539 0.4953 0.25	00 0.2935 0.4478 0.25	25 0.3326 0.4009 0.26	50 0.3701 0.3559 0.27	75 0.4053 0.3136 0.28	00 0.4378 0.2746 0.28	25 0.4672 0.2394 0.29	850 0.4934 0.2079 0.2987	75 0.5164 0.1803 0.30	00 0.5365 0.1561 0.30	25 0.0003 0.7997 0.20	50 0.5539 0.1353 0.31	75 0.5689 0.1173 0.31	
_		.200	.200	.200	.200	.201	.202	.203	.205	.208	.212	.216	.222	.228	.235	.243	.251	.259	.266	.274	.281	.288	.293	0.2990	.303	.307	.311	.314	.316	
580.0 ATM	NH3	.799	.799	. 798	.795	. 792	.786	.777	.765	.748	.726	.698	.665	.627	.585	.540	.493	.445	.398	.353	.311	.272	.237	0.2057	.178	.154	.133	.115	.100	
SSURE = 5	H2	.000	.000	.001	.003	-006	.011	.018	.028	.043	.061	.084	.111	.143	.178	.216	255	.295	.334	.372	.407	.439	.469	0.4952	.518	.538	. 555	.570	.582	
PRESSI	Σ	0	2	5	-	0	N	5	-	0	~	3	~	\circ	N	5	-	0	N	S		0	2	850	-	0	N	5	-	
	N2	.200	.200	.200	. 200	.201	.202	. 203	.205	.208	-212	-217	-222	.229	.236	.243	.251	.259	.267	.274	.281	.288	.294	0.2994	.304	.307	.311	.314	.316	1
570-0 ATM.	NH3	_	_	-	-		-	-	_	_	-	O	O	VQ.	S	M)	J	-4	43.1	(43)	n	N	\sim	0,2034	-	-	-	-	O	
u	2				•																			0.4971					્, •	
PRESSURE	TEMP	300	325	350	375	400	425	450	475	200	525	550	575	009	625	650	675	700	725	750	775	800	825	850	875	006	925	950	975	

PRESS	SSURE = 6	600.0 ATM.	•	PRESSURE	9	10.0 ATM.	•
TEMP	H2	NH3		TEMP	H2	NH3	NZ
300	.59	.08	.31	300	000	_	.200
325	00.	67.	.20	325	.000	-	2
350	.00	19	.20	350	.001	-	.200
375	0.0033	0.7960	0.2007	375	0.0033	0.7961	0.2006
400	.00	.79	.20	400	.006		.201
425	.01	.78	.20	425	.010		- 202
450	.01	.77	.20	450	.017	-	.203
475	.02	.76	.20	475	.027		. 205
500	.04	.74	.20	500	.041	-	.208
525	.06	.72	.21	525	.059	-	.211
550	.08	.70	.21	550	.081		.216
575	.10	•66	.22	575	.108	.0	.221
900	.14	.63	.22	9009	.139	V	.227
625	.17	.58	•23	625	.174	K)	.234
650	.21	.54	.24	650	.211	S	.242
675	.25	49	.25	675	.250	v.	.250
700	.29	.45	.25	700	.289	·5	.257
725	.33	40	.26	725	.328	*	.265
750	.36	.35	-27	750	.365	(L)	.273
775	40	.31	.28	775	.40I	m	.280
800	.43	.27	.28	800	.433	S	.286
825	• 46	.24	*29	825	.463	N	.292
850	.49	.21	• 29	850	.489	N	.297
875	.51	.18	*30	875	.513	, -	.302
900	. 53	.15	.30	006	.533	-	.306
925	.55	.13	.31	925	.551	1444	.310
950	.56	111	.31	950	.566	gered.	.313
975	.58	.10	.31	975	.579	-	-315
1000	• 59	60	.31	1000	.590	0	.318